



UNIVERSITÀ DEGLI STUDI DI MILANO  
DIPARTIMENTO DI FISICA



# *On the connection between memory effects and information exchange between system and environment*

Bassano Vacchini

Dipartimento di Fisica  
“Aldo Pontremoli”  
Università degli Studi di Milano

INFN  
Sezione di Milano

Math-Phys seminars

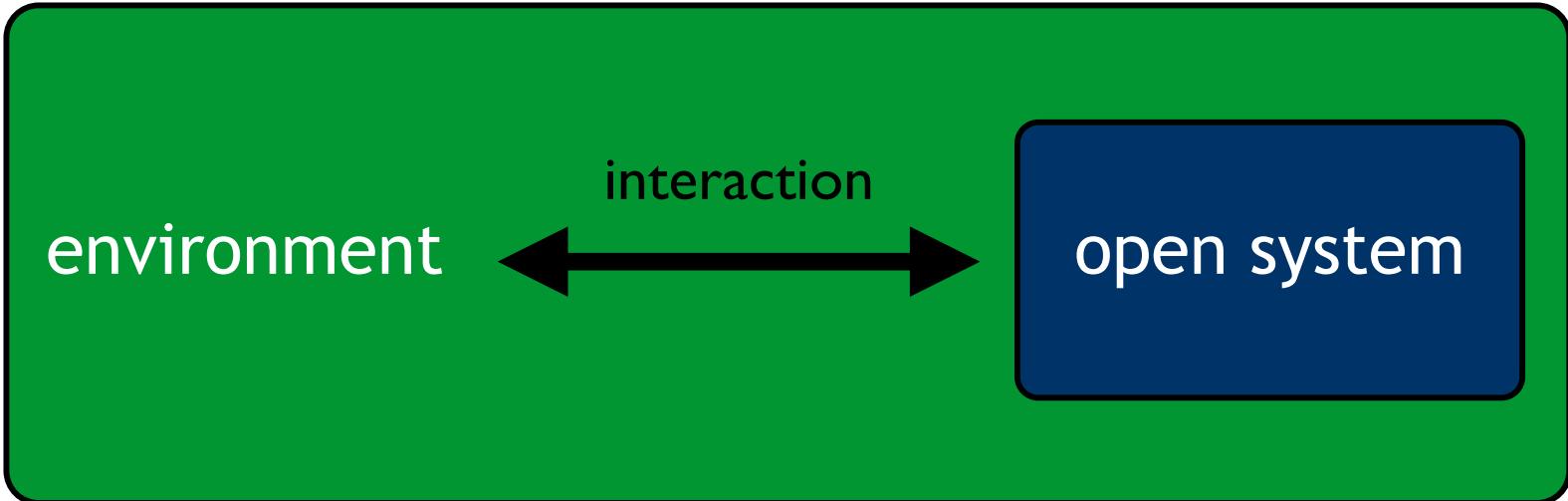
Helsinki (online)  
April 2021

# Outline

---

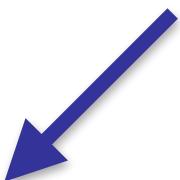
- **Information viewpoint on non-Markovianity in open quantum systems**
- **Entropic bounds on information flow**
- **Reduced vs microscopic dynamics**

# Open quantum systems



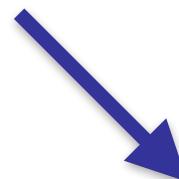
## Bipartite setting

$$H = H_S + H_E + H_I$$
$$\rho_{SE} \in \mathcal{T}(\mathcal{H}_S \otimes \mathcal{H}_E) \quad H \in \mathcal{B}(\mathcal{H}_S \otimes \mathcal{H}_E)$$



## Reduced dynamics

$$\rho_S(0) \rightarrow \rho_S(t) = \Phi(t)\rho_S(0)$$



## Correlations

$$\rho_{SE}(t) \neq \rho_S(t) \otimes \rho_E(t)$$

[Davies, 1976; Alicki & Lendi, 1987; Breuer & Petruccione, 2002; Rivas & Huelga, 2012]

# Quantum process

Time dependent collection of evolution maps

$$\Phi(t)[\rho_S(0)] = \text{Tr}_E(U(t)\rho_S(0) \otimes \rho_E U^\dagger(t)) = \sum_{\alpha,\beta} K_{\alpha,\beta}(t)\rho_S(0)K_{\alpha,\beta}^\dagger(t)$$

emergence of complete positivity

Quantum process

stochasticity of the dynamics  
due to interaction with the environment

on top of

intrinsic probabilistic  
quantum description

[Stinespring PAMS 1955; Hellwig & Kraus CMP 1969; Kraus LNP 1983]



# Memory effects

Obtain & characterize general time evolution law

$$\begin{cases} \frac{d}{dt}\Phi(t) = ? \\ \Phi(t) = ? \end{cases} \quad \text{memoryless or not?}$$

- Connect to and build on classical probability theory
  - classical semigroup vs quantum dynamical semigroup
  - classical Markov process vs quantum Markov process
- Look for physically relevant notion of memory

# Beyond Markovian dynamics\*

## Process viewpoint

$$P_n(t_n, x_n; t_{n-1}, x_{n-1}; \dots, t_1, x_1) \quad t_n \geq t_{n-1} \geq \dots \geq t_1 \geq 0$$

[Lindblad CMP 1979; B. V. & al. NJP 2011; Milz & Modi, arXiv 2020]

## Divisibility viewpoint

$$\Phi(t, \tau)\Phi(\tau, s) = \Phi(t, s) \quad t \geq \tau \geq s \geq 0$$

[Rivas, Huelga & Plenio, PRL 2010; Rivas, Huelga & Plenio, RMP 2014]

## Trajectory viewpoint

$$|\psi(t)\rangle \quad t \geq 0$$

[Piilo & al., PRL 2008; Smirne & al., PRL 2020; Donvil & Muratore-Ginanneschi, arXiv 2021]

## Distinguishability viewpoint

$$D(\rho_1(t), \rho_2(t)) \quad t \geq 0$$

[Breuer, Laine & Piilo, PRL 2009; Breuer, Laine, Piilo & B.V., RMP 2016]

\* Equations and references are a guide for the eye

# Trace distance and distinguishability

## Trace distance

Trace norm natural metric on the space of quantum states

$$D(\rho_1, \rho_2) = \frac{1}{2} \|\rho_1 - \rho_2\| = \frac{1}{2} \text{Tr} |\rho_1 - \rho_2| \quad \|X\| = \sum_i |x_i|$$
$$0 \leq D \leq 1$$

$$D(\rho_1, \rho_2) = 0 \iff \rho_1 = \rho_2$$

$$D(\rho_1, \rho_2) = 1 \iff \rho_1 \perp \rho_2$$

(C)PT maps contractions for the trace distance

$$D(\Phi\rho_1, \Phi\rho_2) \leq D(\rho_1, \rho_2)$$

# Trace distance and distinguishability

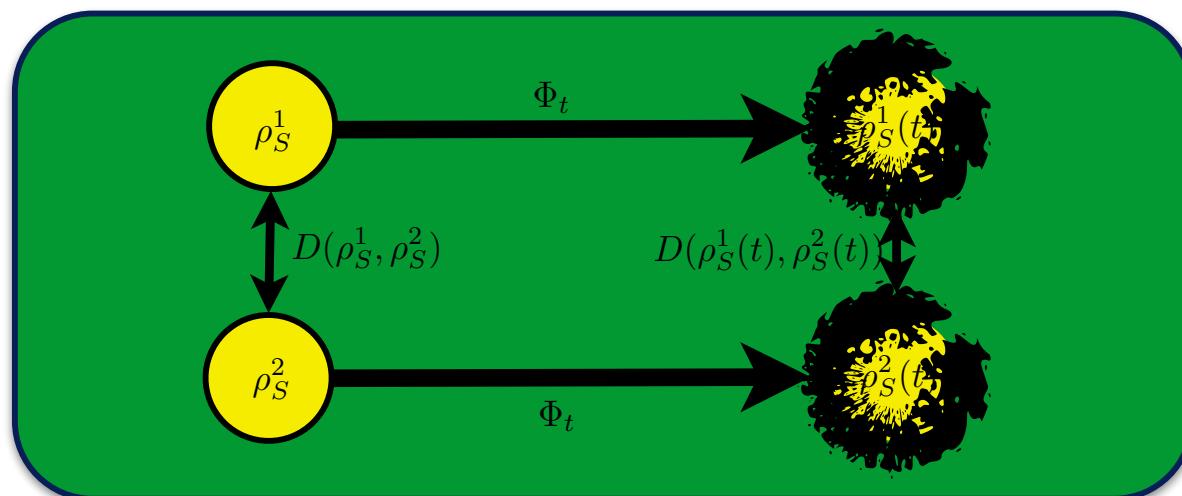
## Distinguishability

Preparations  $\rho_S^1$  and  $\rho_S^2$  taking place with equal frequency to be distinguished upon single measurement

Trace distance determines success of optimal strategy

$$P_{\text{success}} = \frac{1}{2}(1 + D(\rho_S^1, \rho_S^2))$$

Distinguishability decreases under **CPT** map



[Helstrom 1976; Fuchs & al., IEEE 1999]

# Information backflow

## Internal vs external information

$$\mathcal{J}_{int}(t) = D(\rho_S^1(t), \rho_S^2(t)) \quad \mathcal{J}_{ext}(t) = D(\rho_{SE}^1(t), \rho_{SE}^2(t)) - D(\rho_S^1(t), \rho_S^2(t))$$

so that

$$\mathcal{J}_{int}(t) + \mathcal{J}_{ext}(t) = \text{cost}$$

$$\frac{d}{dt} \mathcal{J}_{int}(t) = - \frac{d}{dt} \mathcal{J}_{ext}(t)$$

Non-Markovian behaviour associated to

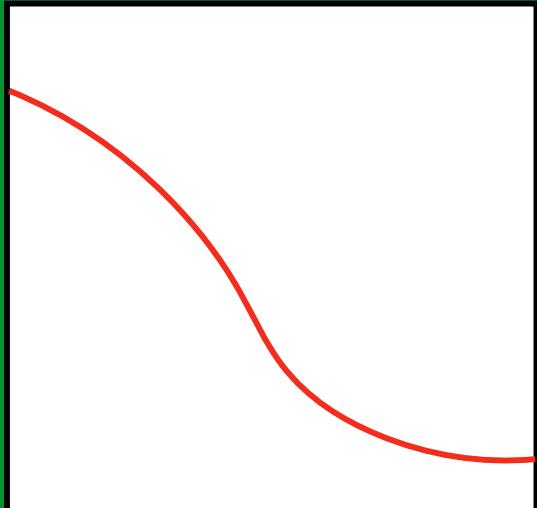
$$\mathcal{J}_{int}(t) - \mathcal{J}_{int}(s) \geq 0 \quad \text{for} \quad t \geq s$$

$$\mathcal{J}_{ext}(t) - \mathcal{J}_{ext}(s) \leq 0 \quad \text{for} \quad t \geq s$$

# Markovian versus non-Markovian dynamics

## Monotonic loss of distinguishability

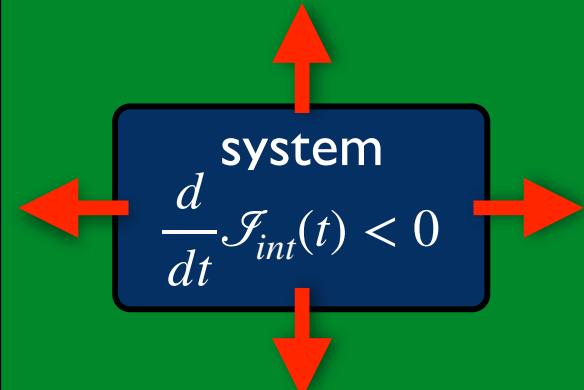
$$D(\rho_S^1(t), \rho_S^2(t))$$



Markovian

$t$

environment



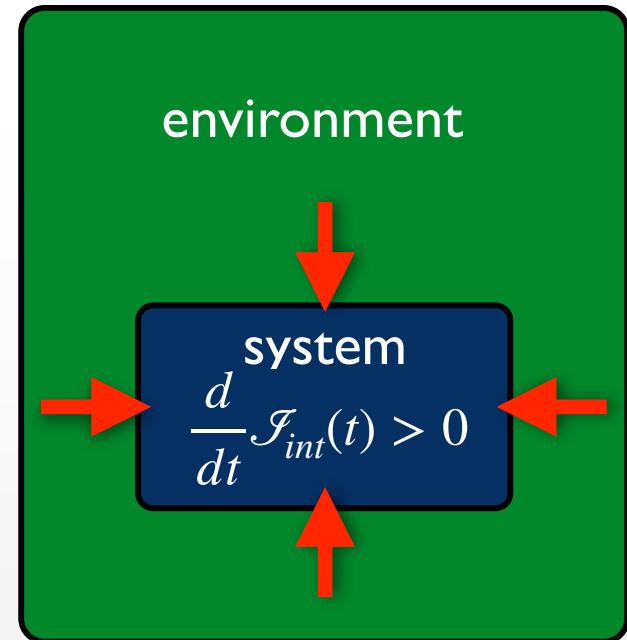
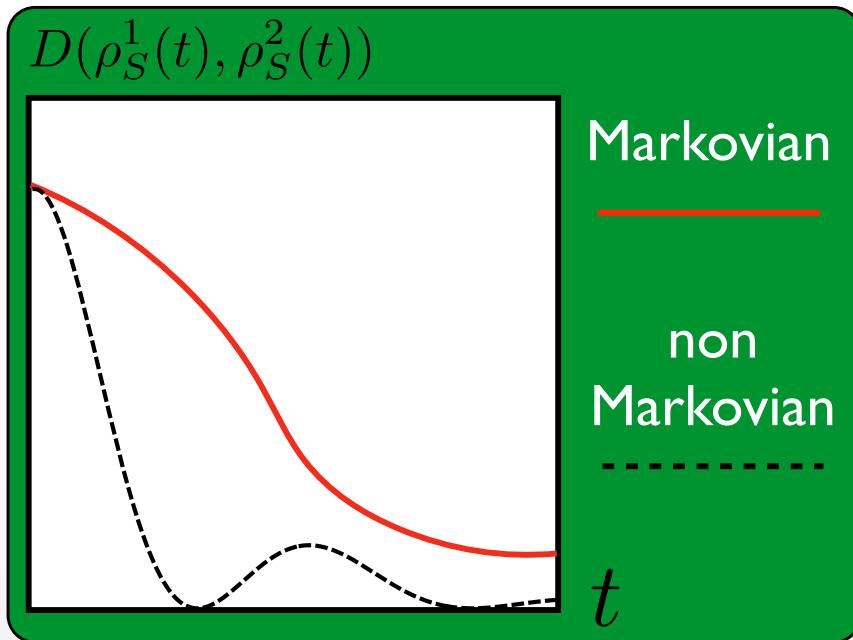
as in the presence of a well-defined composition law

$$D(\rho_1(t), \rho_2(t)) \leq D(\rho_1(s), \rho_2(s)) \quad \forall t \geq s \quad \forall \rho_1(0), \rho_2(0) \in \mathcal{S}(\mathcal{H})$$

Dynamics is said to be Markovian

# Markovian versus non-Markovian dynamics

## Revival of distinguishability



e.g. due to revival in physical property

$$\exists \rho_1(0), \rho_2(0) \in \mathcal{S}(\mathcal{H}) \quad \exists t \geq s \quad D(\rho_1(t), \rho_2(t)) > D(\rho_1(s), \rho_2(s))$$

Dynamics is said to be non-Markovian

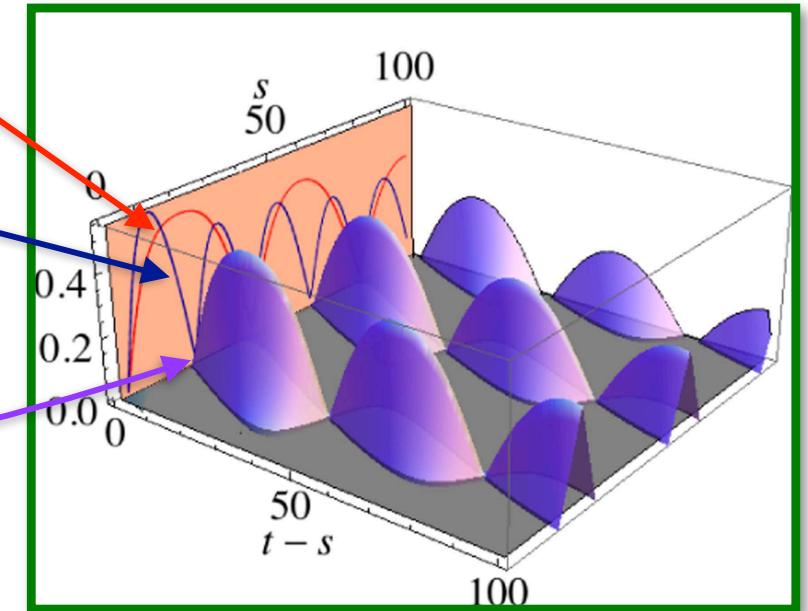
# Information backflow

## Internal vs external information

$$\mathcal{J}_{int}(t) - \mathcal{J}_{int}(s) = \mathcal{J}_{ext}(s) - \mathcal{J}_{ext}(t) \leq \mathcal{J}_{ext}(s)$$

leads to bound on revivals based on external information at previous times comparing states with product of their marginals and comparing environmental marginals

$$\begin{aligned} \mathcal{J}_{int}(t) - \mathcal{J}_{int}(s) &\leq D(\rho_E^1(s), \rho_E^2(s)) \\ &+ D(\rho_{SE}^1(s), \rho_S^1(s) \otimes \rho_E^1(s)) \\ &+ D(\rho_{SE}^2(s), \rho_S^2(s) \otimes \rho_E^2(s)) \end{aligned}$$



# Distinguishing states

## Quantum divergence $f$

Divergence as difference quantifier between classical or quantum probability distributions

Relaxing symmetry and triangle inequality

## Relevant property of quantum divergence $f$

Boundedness

$$0 \leq f(\rho, \sigma) \leq \text{cost}$$

Contractivity under (C)PT maps

$$f(\Phi[\rho], \Phi[\sigma]) \leq f(\rho, \sigma) \implies \begin{cases} f(\mathcal{U}[\rho], \mathcal{U}[\sigma]) = f(\rho, \sigma) \\ f(\rho \otimes \eta, \sigma \otimes \eta) = f(\rho, \sigma) \end{cases}$$

Crucially unitary evolution, partial trace, assignment map are (C)PT transformation

# Microscopic interpretation

## Triangle-like inequality

Assume validity of inequalities

$$f(\varrho, \sigma) - f(\varrho, \tau) \leq \phi_R(f(\sigma, \tau)) \leq \phi(f(\sigma, \tau))$$

$$f(\varrho, \sigma) - f(\eta, \sigma) \leq \phi_L(f(\varrho, \eta)) \leq \phi(f(\varrho, \eta))$$

with monotonic subadditive function  $\phi$

Sufficient condition to derive the bound

$$\begin{aligned} f(\varrho_s(t), \sigma_s(t)) - f(\varrho_s(s), \sigma_s(s)) &\leq \phi \circ \phi(f(\varrho_e(s), \sigma_e(s)) + \\ &\quad \phi(f(\varrho(s), \varrho_s(s) \otimes \varrho_e(s))) + \phi(f(\sigma(s), \sigma_s(s) \otimes \sigma_e(s))) \end{aligned}$$

Connecting distinguishability revivals with correlations and environment changes

Quantum divergence  $\rightarrow$  distance

$\phi \rightarrow 1$

# Outline

---

- Information viewpoint on non-Markovianity in open quantum systems
- Entropic bounds on information flow
- Reduced vs microscopic dynamics

# Entropic divergences

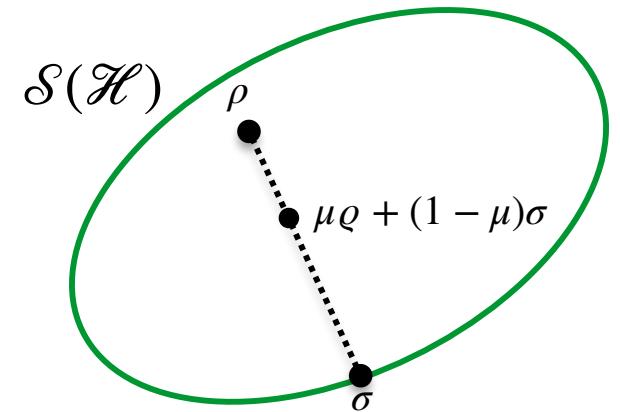
## Telescopic relative entropy

Regularize quantum relative entropy

$$S_\mu(\rho, \sigma) = \frac{1}{\log(1/\mu)} S(\rho, \mu\rho + (1 - \mu)\sigma)$$

$$\mu \in (0,1)$$

$$S(\rho, \sigma) = \text{Tr}\rho \log \rho - \text{Tr}\rho \log \sigma$$



Telescopic relative entropy or quantum skew divergence

Boundedness

$$0 \leq S_\mu(\rho, \sigma) \leq 1$$

Contractivity under (C)PT maps

$$S_\mu(\Phi[\rho], \Phi[\sigma]) \leq S_\mu(\rho, \sigma)$$

# Entropic divergences

## Telescopic relative entropy

Joint convexity

$$\begin{aligned} S_\mu(\lambda\rho_1 + (1-\lambda)\rho_2, \lambda\sigma_1 + (1-\lambda)\sigma_2) \\ \leq \lambda S_\mu(\rho_1, \sigma_1) + (1-\lambda)S_\mu(\rho_2, \sigma_2) \end{aligned}$$

Triangle-like inequalities

$$S_\mu(\varrho, \sigma) - S_\mu(\eta, \sigma) \leq \frac{D(\varrho, \eta)}{\log(1/\mu)} \log \left( 1 + \frac{1}{D(\varrho, \eta)} \frac{1-\mu}{\mu} \right)$$

$$S_\mu(\varrho, \sigma) - S_\mu(\varrho, \tau) \leq \frac{1}{\log(1/\mu)} \log \left( 1 + D(\sigma, \tau) \frac{1-\mu}{\mu} \right)$$

# Entropic divergences

## Telescopic relative entropy

Telescopic Pinsker inequality

$$D(\rho, \sigma) \leq \frac{\sqrt{\log(1/\mu)/2}}{1 - \mu} \sqrt{S_\mu(\rho, \sigma)}$$

together with estimate  $\log(1 + x) \leq \sqrt{x}$

leads to triangle-like inequalities

$$S_\mu(\rho, \sigma) - S_\mu(\rho, \tau) \leq \phi(S_\mu(\sigma, \tau))$$

$$S_\mu(\rho, \sigma) - S_\mu(\eta, \sigma) \leq \phi(S_\mu(\rho, \eta))$$

with

$$\phi(x) = \kappa(\mu) \sqrt[4]{x} \quad \kappa(\mu) = 1 / \sqrt[4]{2\mu^2 \log^3(1/\mu)}$$

[Megier, Smirne & B.V., arXiv 2020]

# Entropic bound on information flow

## Telescopic relative entropy

Straightforward upper bound can be improved to

$$S_\mu(Q_s(t), \sigma_s(t)) - S_\mu(Q_s(s), \sigma_s(s)) \leq \kappa(\mu) \left( \sqrt[4]{S_\mu(Q_E(s), \sigma_E(s))} + \sqrt[4]{S_\mu(Q(s), Q_s(s) \otimes Q_E(s))} + \sqrt[4]{S_\mu(\sigma(s), \sigma_s(s) \otimes \sigma_E(s))} \right)$$

where

$$\kappa(\mu) = 1 / \sqrt[4]{2\mu^2 \log^3(1/\mu)}$$

with minimum value

$$\kappa = (4e^3/27)^{1/4} \approx 1.31 \text{ at } \mu = e^{-3/2}$$

# Entropic divergences

## Symmetrized telescopic relative entropy

Boundedness of telescopic relative entropy makes it natural to symmetrize

$$\bar{S}_\mu(\varrho, \sigma) = \frac{1}{2} \left( S_\mu(\varrho, \sigma) + S_\mu(\sigma, \varrho) \right)$$

Special value  $\mu = 1/2$  recovers quantum Jensen-Shannon divergence

$$\bar{S}_\mu(\varrho, \sigma) \equiv J(\varrho, \sigma) = \frac{1}{2} \left( S \left( \varrho, \frac{\varrho + \sigma}{2} \right) + S \left( \sigma, \frac{\varrho + \sigma}{2} \right) \right)$$

Square root of divergence recently proven to be distance

$$\sqrt{J(\varrho, \sigma)} - \sqrt{J(\varrho, \tau)} \leq \sqrt{J(\sigma, \tau)}$$

$$\sqrt{J(\varrho, \sigma)} - \sqrt{J(\eta, \sigma)} \leq \sqrt{J(\varrho, \eta)}$$

$$\phi(x) = x$$

# Entropic bound on information flow

## Quantum Jensen-Shannon divergence

For the symmetrised case and  $\mu = 1/2$  taking the square root we can further improve using distance property

$$\sqrt{J(\rho_s(t), \sigma_s(t))} - \sqrt{J(\rho_s(s), \sigma_s(s))} \leq \sqrt{J(\rho_E(s), \sigma_E(s))} + \\ \sqrt{J(\rho(s), \rho_s(s) \otimes \rho_E(s))} + \sqrt{J(\sigma(s), \sigma_s(s) \otimes \sigma_E(s))}$$

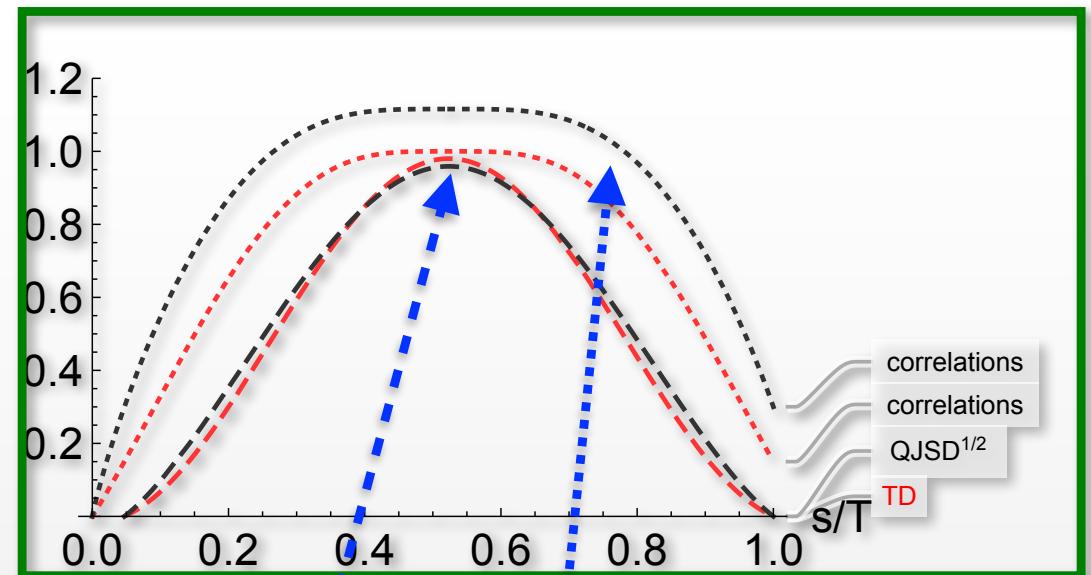
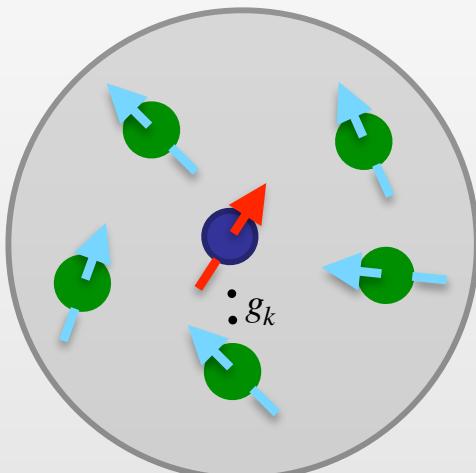
Square root of quantum Jensen-Shannon divergence provides entropy based divergence sharing distance behavior of e.g. trace distance or Bures distance

# Dephasing spin star model

## Revival due to correlations only

Spin star model

$$H_I = \sum_k g_k \sigma_z \otimes \sigma_z^k$$



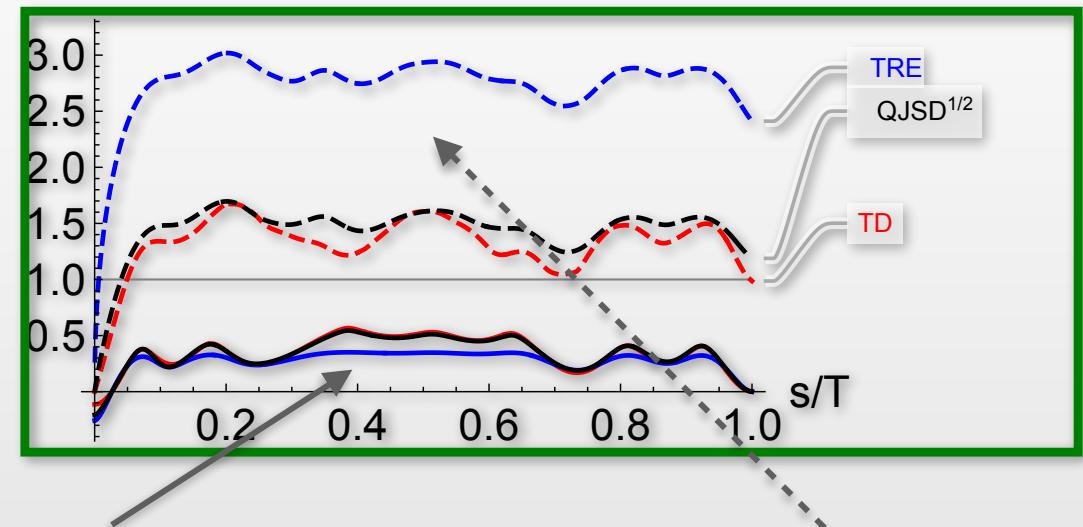
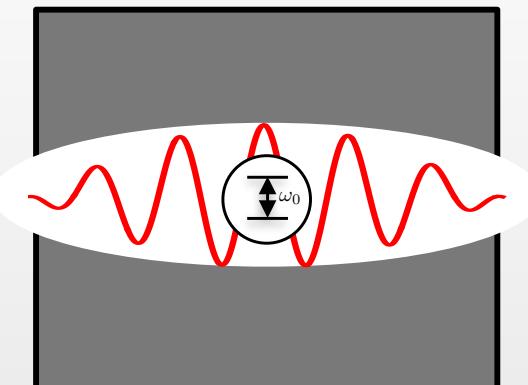
$$\begin{aligned} \mathcal{J}_{int}(t) - \mathcal{J}_{int}(s) &\leq \cancel{f(\rho_E^1(s), \rho_E^2(s))} \\ &+ f(\rho_{SE}^1(s), \rho_S^1(s) \otimes \rho_E^{\cancel{2}}(s)) \\ &+ f(\rho_{SE}^2(s), \rho_S^2(s) \otimes \rho_E^{\cancel{1}}(s)) \end{aligned}$$

# Jaynes-Cummings model

## Divergences and bounds comparison

Two-level system interacting with a bosonic mode

$$H_I = g(\sigma_+ \otimes a + \sigma_- \otimes a^\dagger)$$



$$f(Q_s(T), \sigma_s(T)) - f(Q_s(s), \sigma_s(s)) \leqslant \text{Corr} + \text{Corr} + \text{Env}$$

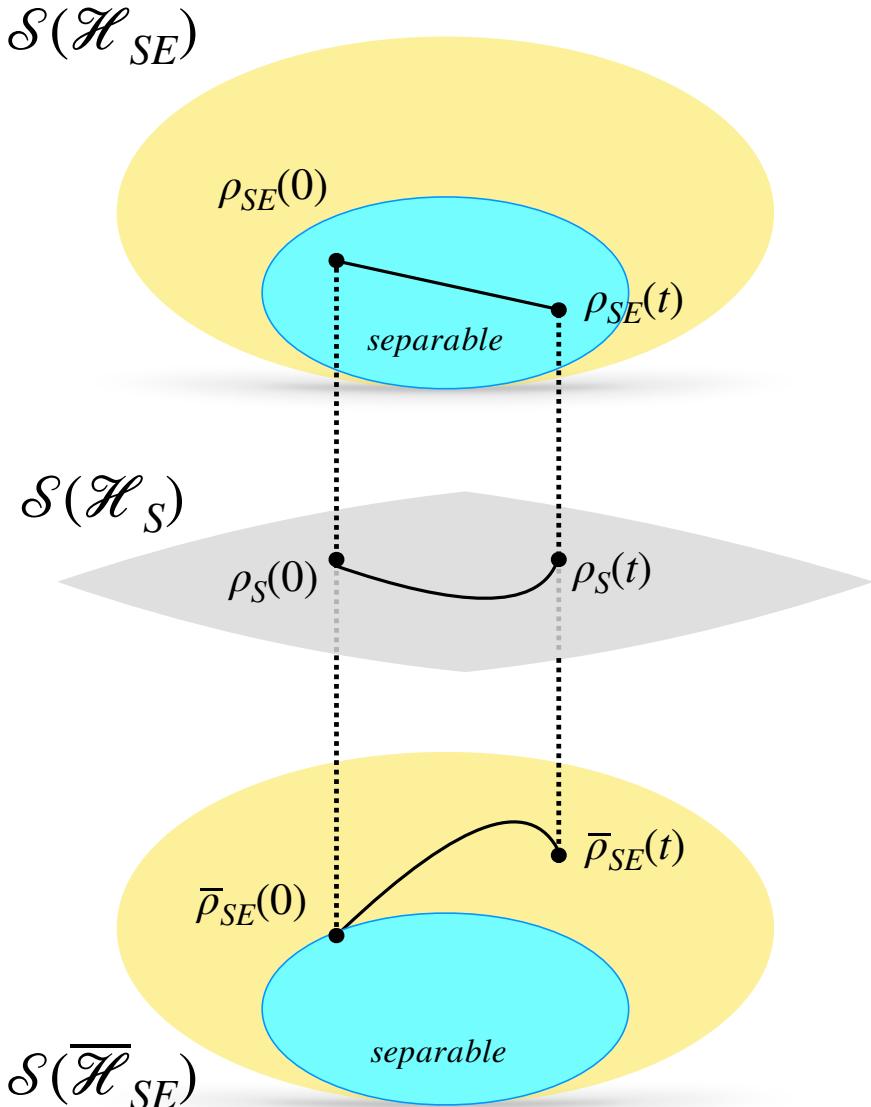
# Outline

---

- **Information viewpoint on non-Markovianity in open quantum systems**
- **Entropic bounds on information flow**
- **Reduced vs microscopic dynamics**

# Reduced vs microscopic description

## Relevance and role of microscopic description



Quantum information viewpoint on quantum non-Markovianity based on local observations

Local signature of microscopic dynamics

# Reduced vs microscopic description

## Different system-environment interaction

Consider system coupled to the same environment initialized in different states, with different coupling terms

$$H_{SE} = H_S + H_E + H_I \quad Q_E(0)$$

$$H_{SE} = H_S + H_E + \bar{H}_I \quad \bar{Q}_E(0)$$

## Same reduced dynamics

Constrain reduced dynamics to be the same for all initial system states

$$\text{Tr}_E\{U_{SE}Q_S(0) \otimes Q_E(0)U_{SE}^\dagger\} = \text{Tr}_E\{\bar{U}_{SE}Q_S(0) \otimes \bar{Q}_E(0)\bar{U}_{SE}^\dagger\}$$
$$\forall Q_S(0) \in \mathcal{S}(\mathcal{H})$$

# Generalized dephasing model

Simplify setting to obtain exact result

$$H_I = \sum_n |n\rangle\langle n| \otimes B_n \quad \bar{H}_I = \sum_n |n\rangle\langle n| \otimes \bar{B}_n$$

Condition for same reduced dynamics becomes

$$\text{Tr}_E e^{-iBt} Q_E(0) = \text{Tr}_E e^{-i\bar{B}t} \bar{Q}_E(0) \iff \text{Tr}_E B^k Q_E(0) = \bar{B}^k \bar{Q}_E(0)$$

Satisfied for the choice

$$\rho_E(0) = \frac{1}{2} (1 + \alpha \cdot \sigma) \quad B = g\eta \cdot \sigma \iff \alpha \cdot \eta = \bar{\alpha} \cdot \bar{\eta}$$

$$\bar{\rho}_E(0) = \frac{1}{2} (1 + \bar{\alpha} \cdot \sigma) \quad \bar{B} = g\bar{\eta} \cdot \sigma$$

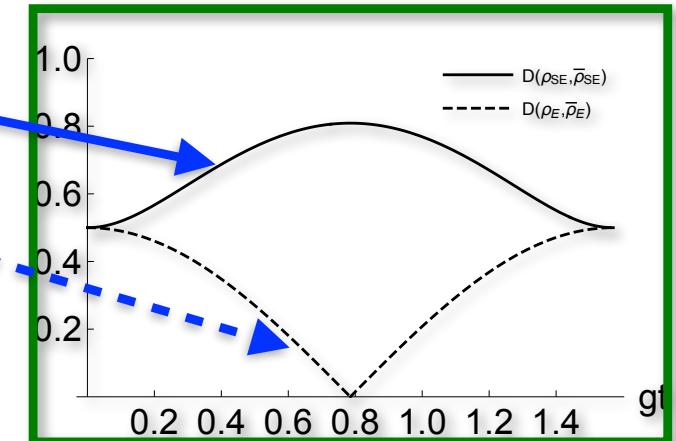
# Generalized dephasing model

Different  $SE$  dynamics

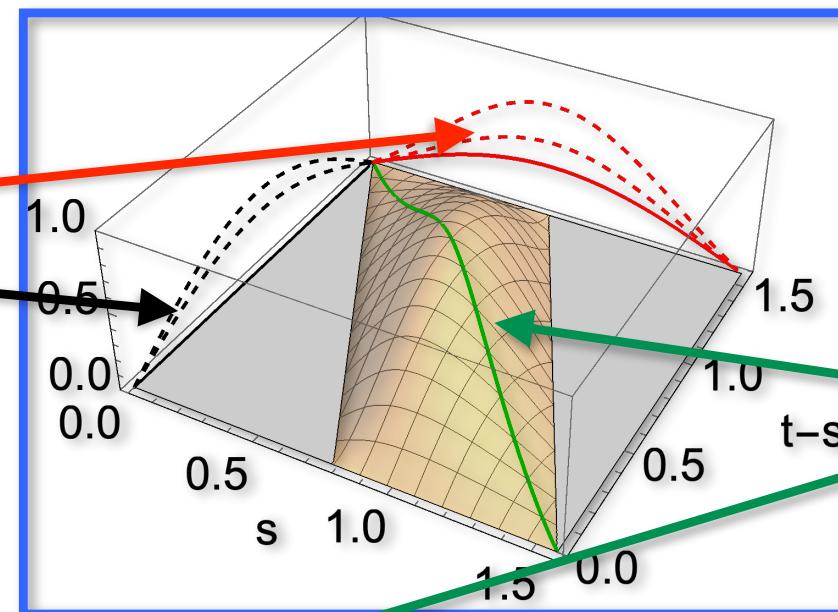
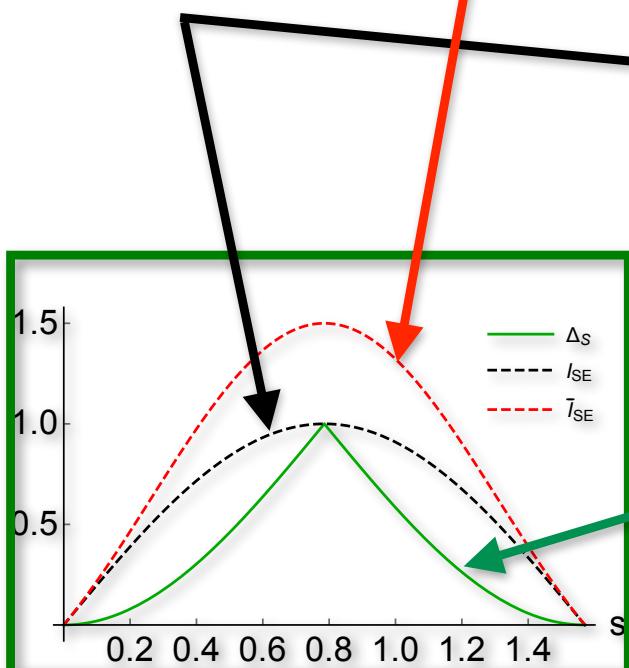
different  $E$  dynamics

different correlations:

- $\rho_{SE}(t)$  remains zero-discord state
- $\bar{\rho}_{SE}(t)$  exhibits entanglement



Different  
external  
information



Identical  
reduced  
dynamics

$$\begin{aligned} \alpha &= (0,0,c) & \eta &= (0,0,1) \\ \bar{\alpha} &= (0,0,1) & \bar{\eta} &= (\sqrt{1-c^2},0,c) \end{aligned}$$

# Recap

---

- Non-Markovianity and quantum info viewpoint
- Divergences of distance and entropic type
- Telescopic entropy and Jensen-Shannon
- Same reduced dynamics with different interactions and bath

N. Megier, A. Smirne and B. Vacchini  
Entropic bounds on information backflow  
[arXiv:2101.02720](https://arxiv.org/abs/2101.02720)

A. Smirne, N. Megier and B. Vacchini  
On the connection between microscopic description and memory effects  
in open quantum system dynamics  
[arXiv:2101.07282](https://arxiv.org/abs/2101.07282) to appear in Quantum

# Acknowledgements

---

Thanks for your attention!

---

**Unimi**  
N. Megier  
A. Smirne

**Collaborations**  
H.-P. Breuer  
S. Campbell



UNIVERSITÀ  
DEGLI STUDI  
DI MILANO

