

KAM stability for conserved quantities in quantum systems

Paolo Facchi - University of Bari & INFN Bari

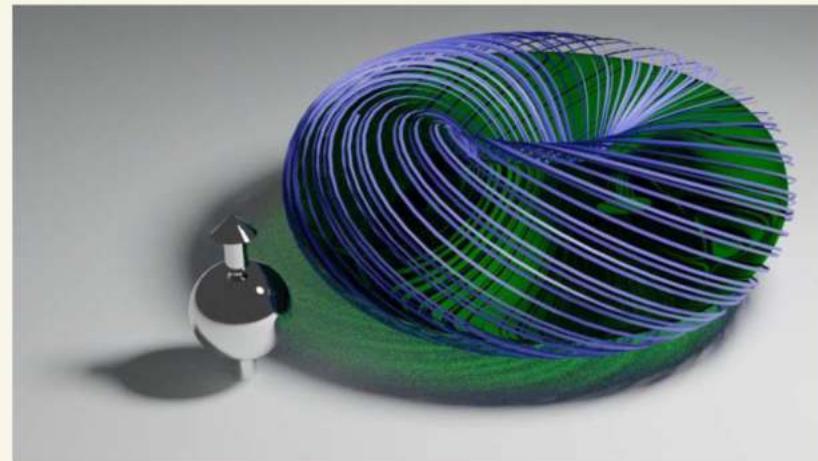
Joint work with

Daniel Burgarth

Hiromichi Nakazato

Fabrizio Pascazio

Kazuya Yusa



Classical mechanics

Kolmogorov - Arnold - Moser (KAM) theorem

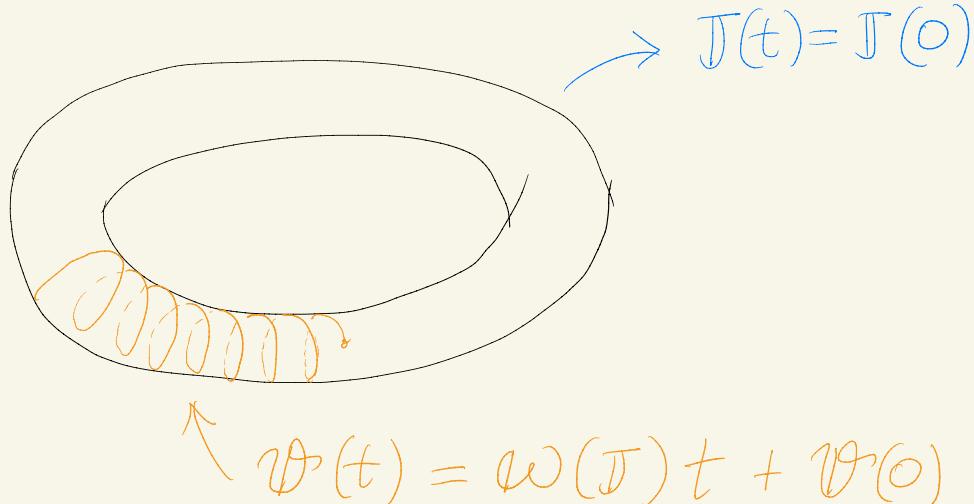
$H_0 = H_0(\mathbb{J})$ integrable system [action-angle variables $(\mathbb{J}, \varPhi) \in \mathbb{R}^d \times \mathbb{T}^d$]

Hamilton equations $\begin{cases} \dot{\varPhi} = \nabla_{\mathbb{J}} H_0(\mathbb{J}) = \omega(\mathbb{J}) \\ \dot{\mathbb{J}} = -\nabla_{\varPhi} H_0(\mathbb{J}) = 0 \end{cases}$

$\Rightarrow \mathbb{J}(t) = \mathbb{J}(0)$ is conserved (actions)

$$\dot{\varPhi}(t) = \varPhi(0) + \omega(\mathbb{J})t \quad (\text{angles})$$

(quasi-)periodic motion on the torus \mathbb{T}^d at $\mathbb{J} = \text{const}$



Consider a small perturbation $\varepsilon \ll 1$

$$H_0(\mathbb{J}) \rightarrow H(\mathbb{J}, \psi) = H_0(\mathbb{J}) + \varepsilon V(\mathbb{J}, \psi)$$

$$\dot{\mathbb{J}}(t) = -\nabla_{\psi} H = \varepsilon \nabla_{\psi} V(\mathbb{J}, \psi) \quad \text{no longer conserved}$$

$$\dot{\psi}(t) = \nabla_{\mathbb{J}} H = \omega(\mathbb{J}) + \varepsilon \nabla_{\mathbb{J}} V(\mathbb{J}, \psi)$$

Theorem (KAM) [1954 - 1967]

Assume H analytic and H_0 nondegenerate $\left[\left(\frac{\partial^2 H_0}{\partial \mathbb{J}_j \partial \mathbb{J}_k}(\mathbb{J}) \right) \text{ invertible} \right]$

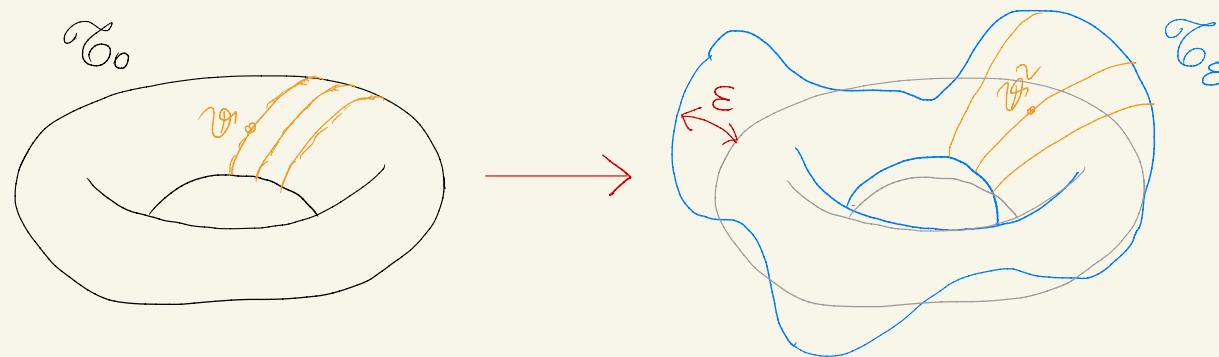
Assume that the torus $T_0 = \{ \mathbb{J}_0 \} \times \mathbb{T}^d$ has frequencies $\omega_0 = \omega(\mathbb{J}_0)$ satisfying

$$|\omega_0 \cdot m| \geq \frac{\gamma}{|m|^\mu} \quad \forall m \in \mathbb{Z}^d \setminus \{0\}, \text{ for some } \gamma > 0, \mu > d-1 \quad (\text{Diophantine condition}).$$

Then $\exists \varepsilon_c > 0$ s.t. $\forall \varepsilon \in (-\varepsilon_c, \varepsilon_c)$, there exists a deformation of the torus T_0 into an invariant torus T_ε .

$$|\omega(\mathbb{J}_0) \cdot m| \geq \frac{\delta}{\|m\|^{\mu}} \quad \forall m \in \mathbb{Z}^d \setminus \{0\}$$

Non-resonant
Diophantine
condition



adiabatic
invariant
eternal

- The motions on T_ε are quasi-periodic with the same frequency ω_0 of the motions on T_0
- The set of Diophantine invariant tori is a Cantor set.
Its complement is dense and has Lebesgue measure $O(\sqrt{\varepsilon})$.

Quantum systems

Conserved quantities vs robustness to perturbations

- fragile symmetries:

small perturbations \rightarrow large deviations

- robust symmetries:

small perturbations \rightarrow small deviations
forever

All symmetries are conserved

but some symmetries are more conserved than others.

Setting

System

Finite-dimensional quantum system (closed)

Hilbert space \mathcal{H} , $\dim \mathcal{H} = d < +\infty$

Hamiltonian

$$H = H^\dagger \in B(\mathcal{H})$$

Evolution

unitary group $t \mapsto e^{-itH}$

Observables

$$A = A^\dagger \in B(\mathcal{H})$$

$$t \mapsto A_t = e^{itH} A e^{-itH}$$

Heisenberg picture

Conserved quantities

$$A_t = e^{itH} A e^{-itH} = A \quad \forall t \in \mathbb{R}$$

\Updownarrow

$$e^{itH} A - A e^{itH} = [e^{itH}, A] = 0 \quad \forall t \in \mathbb{R}$$

\Updownarrow

$$[A, H] = 0$$

Conserved quantities are symmetries [E. Noether]

$$[A, H] = 0 \iff e^{isA} H e^{-isA} = H \quad \forall s \in \mathbb{R}$$

Hamiltonian invariant with respect to the symmetry transformation e^{isA} generated by A .

$Symmetry \iff S = S_t \iff S \in \{H\}^1$

$\{H\}^1 = \{A \in B(\mathcal{H}) : [A, H] = 0\} \quad \text{commutant of } H$

The shape of a symmetry

$$H = \sum_{k=1}^m e_k P_k \quad \text{Spectral resolution} \quad m \leq d$$

$$\text{spec}(H) = \{e_k\} \subset \mathbb{R} \quad \text{spectrum} \quad (e_k \neq e_j, \quad k \neq j)$$

$$P_k P_j = \delta_{kj} P_j, \quad P_k = P_k^+, \quad \sum_{k=1}^m P_k = 1 \quad \text{eigenprojections}$$

$$S \in \{H\}^1 \text{ symmetry iff } S = \sum_{k=1}^m P_k S P_k$$

$$H = \begin{pmatrix} & & \\ & \boxed{e_1} & \\ & e_1 & \\ & & \\ & \boxed{e_2} & e_2 & e_2 \\ & e_2 & e_2 & e_2 \\ & & \boxed{e_3} & \\ & e_3 & e_3 & \\ & & & \end{pmatrix} \Rightarrow S = \begin{pmatrix} & & \\ & \boxed{S_1} & \\ & S_1 & \\ & & \\ & \boxed{S_2} & \\ & S_2 & \\ & & \\ & & \boxed{S_3} \\ & S_3 & \\ & & \end{pmatrix} \quad \text{symmetry}$$

$S \in \{H\}^1$ symmetry iff $S = \sum_{k=1}^m P_k S P_k$

$$[A, H] = 0 \Rightarrow P_k [A, H] P_j = 0$$

$$0 = P_k (AH - HA) P_j = P_k \underbrace{AH P_j}_{\ell_j P_j} - P_k \underbrace{HA P_j}_{\ell_k P_k} = (\ell_j - \ell_k) P_k A P_j$$

$$\Rightarrow P_k A P_j = 0 \text{ if } k \neq j$$

$$\Rightarrow A = \sum_k P_k A \sum_j P_j = \sum_{k,j} P_k A P_j = \sum_k P_k A P_k$$

$$H = \begin{pmatrix} & & \\ & \boxed{\ell_1 \ell_1} & \\ & & \boxed{\ell_2 \ell_2 \ell_2 \ell_2} \\ & & & \boxed{\ell_3 \ell_3} \end{pmatrix} \Rightarrow S = \begin{pmatrix} & & \\ & \boxed{S_1} & \\ & & \boxed{S_2} \\ & & & \boxed{S_3} \end{pmatrix}$$

symmetry

Zeno projection

$$S \in \{H\}^{\perp} \Leftrightarrow S = \sum_{k=1}^m P_k S P_k \quad (\text{it depends only on } P_k \text{ not on } e_k)$$

$$\begin{aligned} \pi: B(\mathcal{H}) &\longrightarrow \{H\}^{\perp} \\ A &\longmapsto \pi(A) = \sum_{k=1}^m P_k A P_k \end{aligned} \quad \begin{array}{l} \text{projection onto} \\ \text{the commutant} \end{array}$$

$$S \in \{H\}^{\perp} \Leftrightarrow S = \pi(S)$$

$$V_Z = \pi(V) = \sum_{k=1}^m P_k V P_k \quad \text{"Zeno" Hamiltonian}$$

$$V_Z = \pi(V) = \sum_{k=1}^m P_k V P_k \quad \text{"Zeno" Hamiltonian}$$

$$\| e^{-it(KH+V)} - e^{-it(KH+V_Z)} \| \leq 2 \frac{\sqrt{m}}{\eta} \frac{1}{K} \|V\| (1 + \|V\|t)$$

$$e^{-it(KH+V)} = e^{-it(KH+V_Z)} + O\left(\frac{1}{K}\right) \quad \text{up to times } O(1)$$

as $K \rightarrow \infty$

strong coupling limit

$$(P e^{-it\frac{V}{n}})^n = \underbrace{P e^{-it\frac{V}{n}} P e^{-it\frac{V}{n}} \dots P e^{-it\frac{V}{n}}}_{n \text{ times}} \rightarrow e^{-itPV} P$$

as $n \rightarrow \infty$

quantum Zeno effect (frequent projective measurements)

What happens if

$$H \rightarrow H + \varepsilon V \quad \varepsilon \ll 1 \quad (\|V\|=1)$$

small perturbation ?

S symmetry, $S \in H^3$?

$$S_t^\varepsilon = e^{it(H+\varepsilon V)} S e^{-it(H+\varepsilon V)} \stackrel{?}{=} S + \overset{\text{independent}}{\underset{\text{of } t}{O(\varepsilon)}} \quad \text{for all } \|V\|=1$$

By perturbation theory this is true for small times. Is it stable for all times?

If this is the case, we call S a robust symmetry.

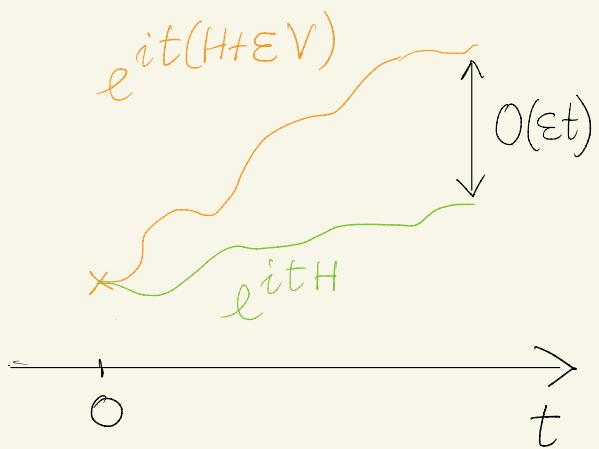
Divergence of two quantum evolutions

$$\begin{aligned}
 e^{it(H+\varepsilon V)} - e^{itH} &= \left[e^{i(t-s)H} e^{is(H+\varepsilon V)} \right]_{s=0}^{s=t} \\
 &= \int_0^t ds \frac{d}{ds} \left[e^{i(t-s)H} e^{is(H+\varepsilon V)} \right] \\
 &= \int_0^t ds e^{i(t-s)H} i\varepsilon V e^{is(H+\varepsilon V)}
 \end{aligned}$$

$$\|e^{it(H+\varepsilon V)} - e^{itH}\| \leq \int_0^t ds \varepsilon \|e^{i(t-s)H} V e^{is(H+\varepsilon V)}\|$$

$$\|WAU\| = \|A\| \quad (*) \quad = \int_0^t ds \varepsilon \|V\| = \varepsilon t \|V\|$$

$$\delta(t) = \|e^{it(H+\varepsilon V)} - e^{itH}\| \leq \varepsilon \|V\| t$$



$$(*) \quad \|WAU\| = \|A\| \quad \forall W, U \text{ unitaries}$$

$$\begin{aligned}
 \|WAU\| &= \sup_{\|\varphi\|=\|\psi\|=1} |\langle \varphi | WAU \psi \rangle| \\
 &= \sup_{\|\varphi\|=\|\psi\|=1} |\langle W^* \varphi | A \underbrace{U \psi}_{\tilde{\psi}} \rangle| \\
 &= \sup_{\|\tilde{\psi}\|=\|\psi\|=1} |\langle \tilde{\psi} | A \tilde{\psi} \rangle| = \|A\|
 \end{aligned}$$

Evolution of observables

$$\|A_t^\varepsilon - A_t\| \leq 2\|A\| \|e^{it(H+\varepsilon V)} - e^{itH}\| = 2\|A\| \delta(t)$$

$$\begin{aligned}
 \|A_t^\varepsilon - A_t\| &= \|e^{it(H+\varepsilon V)} A e^{-it(H+\varepsilon V)} - e^{itH} A e^{-itH}\| \\
 &= \|e^{itH} (e^{-itH} e^{it(H+\varepsilon V)} A - A e^{-itH} e^{it(H+\varepsilon V)}) e^{-it(H+\varepsilon V)}\| \\
 \|WAU\| &= \|A\| \quad (*) \\
 &= \|[e^{-itH} e^{it(H+\varepsilon V)}, A]\| \\
 &= \|[e^{-itH} e^{it(H+\varepsilon V)} - 1, A]\| \quad [1, A] = 0 \\
 &\leq 2\|A\| \|e^{-itH} e^{it(H+\varepsilon V)} - 1\| \quad (*) = 2\|A\| \|e^{it(H+\varepsilon V)} - e^{itH}\|
 \end{aligned}$$

$$\|A_t^\varepsilon - A_t\| \leq 2\|A\| \delta(t) \leq 2\varepsilon t \|V\| \|A\| = O(\varepsilon) \quad \text{for } t = O(1)$$

For a symmetry : $S_t = e^{itH} S e^{-itH} = S$

$$\|S_t^\varepsilon - S\| \leq 2\varepsilon t \|\nabla\| \|S\| = O(\varepsilon) \quad \text{for } t = O(1)$$

However, notice that

$$\|S_t^\varepsilon - S\| \leq \|S_t^\varepsilon\| + \|S\| = 2\|S\|$$

maximum distance = $2\|S\|$ obtained if $S_t^\varepsilon = -S$

A symmetry $S \in \mathcal{H}^{\mathcal{G}}_+$ is **fragile** if there exists a perturbation $V = V^* \in B(\mathcal{H})$ such that

$$\sup_{t \in \mathbb{R}} \|e^{+it(H+\varepsilon V)} S e^{-it(H+\varepsilon V)} - S\| = O(1) \quad \forall \varepsilon > 0 \text{ however small}$$

Eventually, for sufficiently long times, it becomes $O(1)$

$$\|S_t^\varepsilon - S\| = O(\varepsilon) \quad \text{for times } t = O(1)$$

$$= O(1) \quad \text{for times } t = O\left(\frac{1}{\varepsilon}\right)$$

the divergence can incrementally accumulate over time

but does it?

$$H|\psi_1\rangle = \epsilon |\psi_1\rangle, \quad H|\psi_2\rangle = \epsilon |\psi_2\rangle$$

$$H = \begin{pmatrix} & & \\ & \boxed{\epsilon_1} & \\ & \epsilon_1 & \\ & & \\ & e & e \\ & e & e \\ & e & \\ & & \\ & \boxed{\epsilon_3} & \\ & \epsilon_3 & \\ & & \end{pmatrix}$$

$$S = \begin{pmatrix} & & \\ & S_1 & \\ & & \\ & \boxed{\delta_1} & \times \\ & \delta_2 & \times \\ & \times & \times & \times \\ & & & \\ & & & S_3 \end{pmatrix}$$

*S breaks
the degeneracy*

$$S|\psi_1\rangle = s_1|\psi_1\rangle, \quad S|\psi_2\rangle = s_2|\psi_2\rangle \quad s_2 - s_1 > 0$$

$$V = |\psi_1\rangle\langle\psi_2| + |\psi_2\rangle\langle\psi_1|$$

$$V = \begin{pmatrix} & & \\ & & \\ & & \\ & \boxed{1} & \\ & 1 & \\ & & \\ & & \\ & & \end{pmatrix}$$

$$|\Psi\rangle = \alpha|\psi_1\rangle + \beta|\psi_2\rangle$$

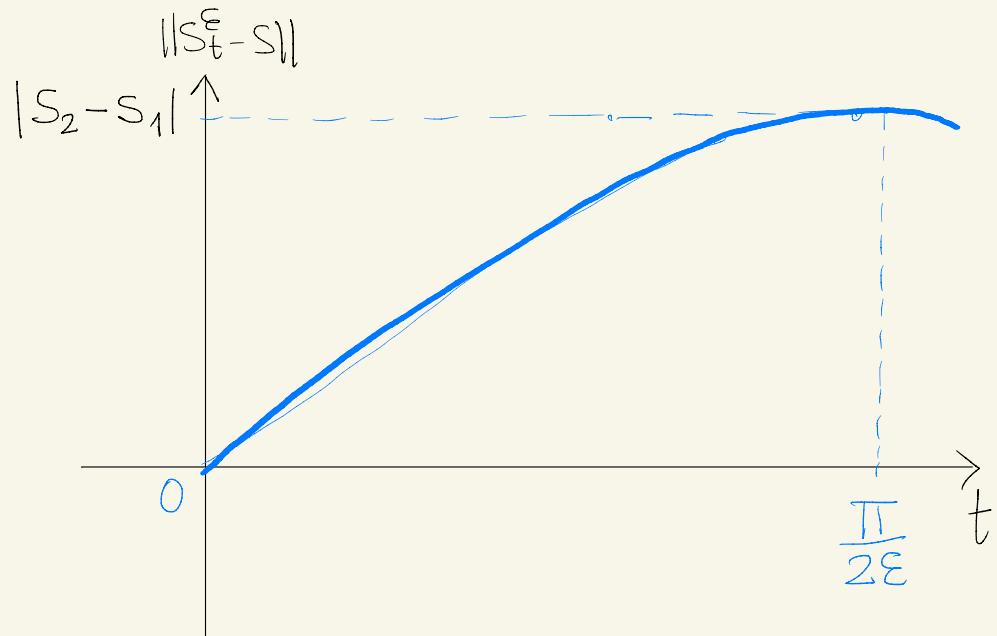
two-dimensional problem

$$|\psi_1(t)\rangle = e^{-it(H+\epsilon V)} |\psi_1\rangle = e^{-ite} (\cos(\epsilon t) |\psi_1\rangle - i \sin(\epsilon t) |\psi_2\rangle)$$

$$|\psi_2(t)\rangle = e^{-it(H+\epsilon V)} |\psi_2\rangle = e^{-ite} (-i \sin(\epsilon t) |\psi_1\rangle + \cos(\epsilon t) |\psi_2\rangle)$$

$$\|S_t^\epsilon - S\| = |S_2 - S_1| |\sin \epsilon t|$$

$$t = \frac{\pi}{2\epsilon} = O\left(\frac{1}{\epsilon}\right) \Rightarrow \|S_t^\epsilon - S\| = O(1)$$



S is fragile over long times!

$$|\psi_1(t)\rangle = e^{-it(H+\epsilon V)} |\psi_1\rangle = e^{-ite} (\cos(\epsilon t) |\psi_1\rangle - i \sin(\epsilon t) |\psi_2\rangle)$$

$$|\psi_2(t)\rangle = e^{-it(H+\epsilon V)} |\psi_2\rangle = e^{-ite} (-i \sin(\epsilon t) |\psi_1\rangle + \cos(\epsilon t) |\psi_2\rangle)$$

$$S_t^\epsilon = S_1 |\psi_1(-t)\rangle \langle \psi_1(-t)| + S_2 |\psi_2(-t)\rangle \langle \psi_2(-t)|$$

$$S_t^\epsilon = \begin{pmatrix} S_1 \cos^2(\epsilon t) + S_2 \sin^2(\epsilon t) & i(S_1 - S_2) \sin(\epsilon t) \cos(\epsilon t) \\ -i(S_1 - S_2) \sin(\epsilon t) \cos(\epsilon t) & S_1 \sin^2(\epsilon t) + S_2 \cos^2(\epsilon t) \end{pmatrix}$$

$$S_t^\epsilon - S = \begin{pmatrix} (S_2 - S_1) \sin^2(\epsilon t) & -i(S_2 - S_1) \sin(\epsilon t) \cos(\epsilon t) \\ i(S_2 - S_1) \sin(\epsilon t) \cos(\epsilon t) & -(S_2 - S_1) \sin^2(\epsilon t) \end{pmatrix}$$

$$S_t^\epsilon - S = (S_2 - S_1) \sin(\epsilon t) \begin{pmatrix} \sin(\epsilon t) & -i \cos(\epsilon t) \\ i \cos(\epsilon t) & -\sin(\epsilon t) \end{pmatrix} = (S_2 - S_1) \sin(\epsilon t) (\cos(\epsilon t) \sigma_y + \sin(\epsilon t) \sigma_z)$$

$$\|\cos(\epsilon t) \sigma_y + \sin(\epsilon t) \sigma_z\| = 1$$

$$\|S_t^\epsilon - S\| = |S_2 - S_1| |\sin(\epsilon t)| \leq 2 \|S\| |\sin(\epsilon t)|$$

Bicommutant $\{H\}''$

$$S = \sum_{k=1}^m s_k P_k \quad S = f(H) \in \{H\}'' \subset \{H\}'$$

$$\{H\}'' = \{ A \in B(\mathcal{H}) : [A, B] = 0 \quad \forall B \in \{H\}' \}$$

$$= \{ f(H) \in B(\mathcal{H}) : f: \mathbb{R} \rightarrow \mathbb{C} \}$$

$$= \{ \sum_{k=1}^m f(e_k) P_k : f: \mathbb{R} \rightarrow \mathbb{C} \}$$

$$= \{ \sum_{k=1}^{m-1} P_{m-1}(e_k) P_k : P_{m-1}(x) = \sum_{k=0}^{m-1} a_k x^k, a_k \in \mathbb{C} \}$$

$$= \{ \sum_{k=0}^m a_k H^k : a_k \in \mathbb{C} \}$$

$S \in \{H\}''$ is robust?

$$(f(e_1), \dots, f(e_m)) \leftrightarrow (a_0, a_1, \dots, a_{m-1})$$

$$\begin{cases} a_0 + a_1 e_1 + a_2 e_1^2 + \dots + a_{m-1} e_1^{m-1} = f(e_1) \\ a_0 + a_1 e_m + a_2 e_m^2 + \dots + a_{m-1} e_m^{m-1} = f(e_m) \end{cases}$$

$$V = \begin{pmatrix} 1 & e_1 & e_1^2 & \dots & e_1^{m-1} \\ 1 & e_2 & e_2^2 & \dots & e_2^{m-1} \\ \vdots & \ddots & \ddots & \ddots & \vdots \\ 1 & e_m & e_m^2 & \dots & e_m^{m-1} \end{pmatrix} = (e_j^{k-1})_{1 \leq j, k \leq m}$$

$$\det V = \prod_{1 \leq j < k \leq m} (e_j - e_k) \neq 0$$

Vandermonde determinant

What if the perturbation is a symmetry?

$$V = \pi(V) \in \{H\}^I$$

$$S = f(H) \in \{H\}^H$$

$$S = \sum_{k=1}^m s_k P_k$$



$$[S, V] = 0$$

$$S_t^\varepsilon = e^{it(H + \varepsilon V)} S e^{-it(H + \varepsilon V)} = S = S_t = e^{ith} S e^{-ith}$$

$$S_t^\varepsilon = S_t = S \quad \text{conserved for all times}$$

For a generic $V \notin \{H\}^\perp$ it is no longer true that $S_t^\varepsilon = S$.

However, we have

$$\begin{aligned}\|S_t^\varepsilon - S\| &= \| [e^{it(H+\varepsilon V)}, S] \| \\ &= \| [e^{it(H+\varepsilon V)} - e^{it(H+\varepsilon \pi(V))}, S] + \underbrace{[e^{it(H+\varepsilon \pi(V))}, S]}_{=0} \| \end{aligned}$$

$$S \in \{H\}^\perp \quad H, \pi(V) \in \{H\}^\perp \Rightarrow [e^{it(H+\varepsilon \pi(V))}, S] = 0$$

$$\|S_t^\varepsilon - S\| \leq 2\|S\| \|e^{it(H+\varepsilon V)} - e^{it(H+\varepsilon \pi(V))}\| = 2\|S\| \delta_Z(t)$$

$$S \in \{H\}^{\parallel}$$

$$\|S_t^\varepsilon - S\| \leq 2\|S\| \delta_Z(t)$$

$$\delta_Z(t) = \|e^{it(H+\varepsilon V)} - e^{it(H+\varepsilon \pi(V))}\|$$

One can prove that $\delta_Z(t) \leq \frac{2\sqrt{m}}{\eta} \varepsilon \|V\| (1 + \varepsilon \|V\| t)$

where $\eta = \min_{k \neq j} |\ell_k - \ell_j|$ minimal spectral gap of H

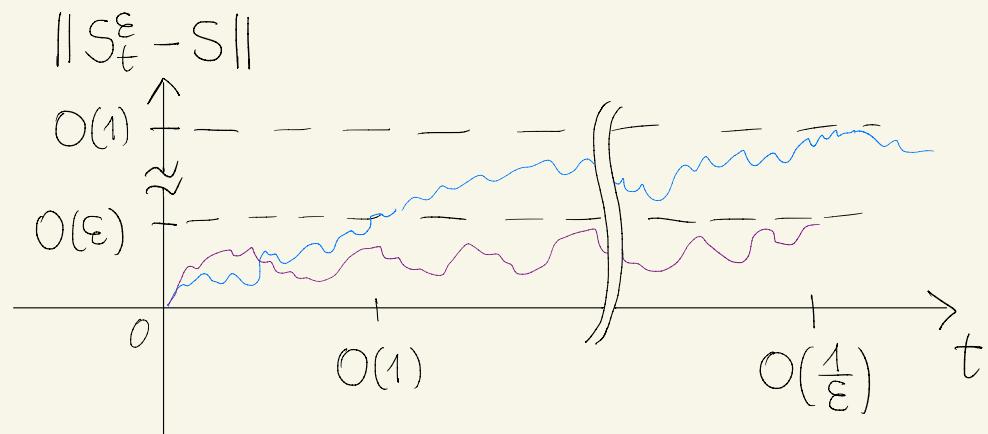
$$\delta_z(t) = O(\varepsilon)(1 + O(\varepsilon t))$$

The distance can accumulate over time.

Notice, however, that

$$\|S_t^\varepsilon - S\| = O(\varepsilon) \quad \text{up to times } t = O\left(\frac{1}{\varepsilon}\right) !$$

The symmetry S is conserved for longer times.



$$S \in \mathcal{H}^1 \setminus \mathcal{H}^{\text{II}}$$

$$S \in \mathcal{H}^{\text{II}}$$

Improved strategy : generic $\hat{V} \in \mathcal{H}\mathcal{H}^{\dagger}$ $\hat{V} \neq \pi(V)$

$$\|S_t^\varepsilon - S\| = \| [e^{it(H+\varepsilon V)}, S] \|$$

$$= \| [e^{it(H+\varepsilon V)} - e^{it(H+\varepsilon \hat{V})}, S] + \underbrace{[e^{it(H+\varepsilon \hat{V})}, S]}_{=0} \|$$

$$S \in \mathcal{H}\mathcal{H}^{\dagger} \quad H, \hat{V} \in \mathcal{H}\mathcal{H}^{\dagger} \Rightarrow [e^{it(H+\varepsilon \hat{V})}, S] = 0$$

$$\|S_t^\varepsilon - S\| \leq 2 \|e^{it(H+\varepsilon V)} - e^{it(H+\varepsilon \hat{V})}\| \|S\| \leq 2 \|S\| \delta_\infty$$

$$\delta_\infty = \sup_{t \in \mathbb{R}} \|e^{it(H+\varepsilon V)} - e^{it(H+\varepsilon \hat{V})}\| = O(\varepsilon) ?$$

MAIN IDEA:

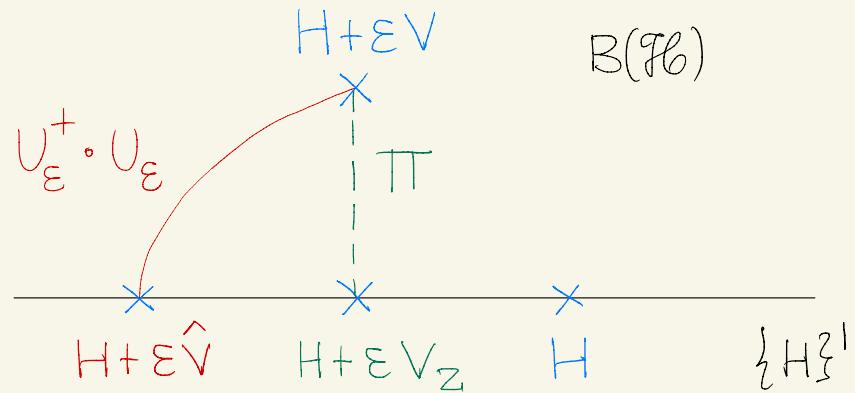
Suppose there exists a unitary U_ε , close to the identity $U_\varepsilon = 1 + O(\varepsilon)$, such that

$$H + \varepsilon \hat{V} = U_\varepsilon^\dagger (H + \varepsilon V) U_\varepsilon$$

isospectral

$$\begin{aligned}
S_{\infty} &= \sup_{t \in \mathbb{R}} \| e^{it(H+\varepsilon V)} - e^{it(H+\varepsilon \hat{V})} \| \\
&= \sup_{t \in \mathbb{R}} \| e^{it(H+\varepsilon V)} - U_\varepsilon^* e^{it(H+\varepsilon V)} U_\varepsilon \| \\
&= \sup_{t \in \mathbb{R}} \| [U_\varepsilon, e^{it(H+\varepsilon V)}] \| \\
&= \sup_{t \in \mathbb{R}} \| [U_\varepsilon - 1, e^{it(H+\varepsilon V)}] \| \\
&\leq 2 \sup_{t \in \mathbb{R}} \| U_\varepsilon - 1 \| \| e^{it(H+\varepsilon V)} \| = 2 \| U_\varepsilon - 1 \| = O(\varepsilon)
\end{aligned}$$

!



$$\text{spec}(H + \varepsilon \hat{V}) = \text{spec}(H + \varepsilon V)$$

Isospectral perturbations

$$H = H^+ \quad \text{perturbation} \quad \tilde{H} = H + O(\varepsilon), \quad \varepsilon \ll 1$$

divergence $\delta(t) = \|e^{it\tilde{H}} - e^{itH}\|$

spectral decomposition : $H = \sum_{k=1}^m \ell_k P_k \quad \tilde{H} = \sum_{k=1}^m \tilde{\ell}_k \tilde{P}_k$

$\tilde{\ell}_k \neq \tilde{\ell}_j \quad \text{for } k \neq j \quad (\text{but it may be } \ell_k = \ell_j)$

$$\tilde{P}_k = P_k + O(\varepsilon) \quad \tilde{\ell}_k = \ell_k + O(\varepsilon) \quad [\text{Kato}]$$

$$\begin{aligned} e^{it\tilde{H}} - e^{itH} &= \sum_k (e^{it\tilde{\ell}_k} \tilde{P}_k - e^{it\ell_k} P_k) \\ &= \underbrace{\sum_k e^{it\tilde{\ell}_k} (\tilde{P}_k - P_k)}_{\text{eigenprojections}} + \underbrace{\sum_k (e^{it\tilde{\ell}_k} - e^{it\ell_k}) P_k}_{\text{eigenvalues}} \end{aligned}$$

$$\delta_P(t) = \left\| \sum_k e^{it\tilde{\ell}_k} (\tilde{P}_k - P_k) \right\| \leq \sum_k \|\tilde{P}_k - P_k\| = O(\varepsilon) \quad \text{uniformly in time}$$

$$\delta_e(t) = \left\| \sum_k (e^{it\tilde{\ell}_k} - e^{it\ell_k}) P_k \right\| = \max_k |e^{it\tilde{\ell}_k} - e^{it\ell_k}| = 2 \max_k \left| \sin \left(t \frac{\tilde{\ell}_k - \ell_k}{2} \right) \right|$$

$$\delta(t) = \delta_e(t) + \delta_P(t) = 2 \max_k \left| \sin \left(t \frac{\tilde{\ell}_k - \ell_k}{2} \right) \right| + O(\varepsilon)$$

$$\delta(t) = O(\varepsilon) \quad \text{for } t = O(1)$$

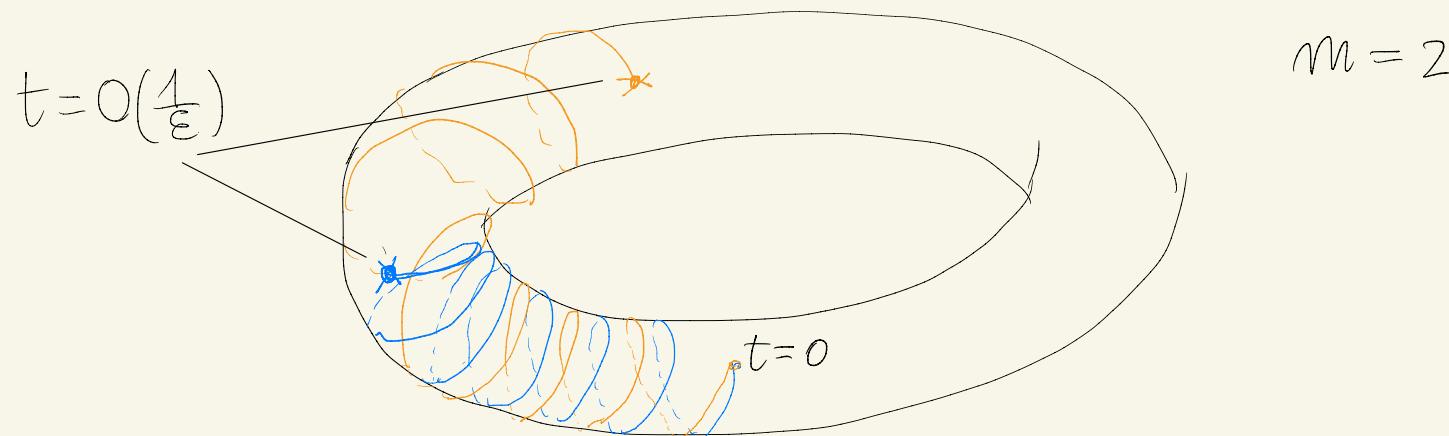
but $\delta(t) = 2 + O(\varepsilon) \quad \text{for } t = \frac{\pi}{\tilde{\ell}_k - \ell_k} = O\left(\frac{1}{\varepsilon}\right)$

$\delta_\ell(t)$ secular drift.

$$\delta_\infty = \sup_t \delta(t) = 2$$

$$\boxed{\delta_\infty = O(\varepsilon) \quad \text{iff} \quad \tilde{\ell}_k = \ell_k \quad \forall k \\ \text{iff} \quad \text{spec}(H) = \text{spec}(\tilde{H})}$$

$$\sum_k (e^{it\tilde{\ell}_k} - e^{it\ell_k}) P_k \xrightarrow[m-\text{torus}]{} (e^{it\tilde{\ell}_1}, e^{it\tilde{\ell}_2}, \dots, e^{it\tilde{\ell}_m}) - (e^{it\ell_1}, e^{it\ell_2}, \dots, e^{it\ell_m})$$



KAM iteration scheme

$$H + \varepsilon \hat{V}(\varepsilon) = U_\varepsilon^+ (H + \varepsilon V) U_\varepsilon \quad \text{with} \quad \hat{V}(\varepsilon) = \Pi(\hat{V}(\varepsilon))$$

$$\Pi(A) = \sum_{k=1}^m P_k A P_k$$

$$U_\varepsilon = e^{i K(\varepsilon)} \quad K(\varepsilon) = \varepsilon K_1 + O(\varepsilon^2) \quad K_1 = K_1^+$$

$$\hat{V}(\varepsilon) = V_0 + O(\varepsilon) \quad V_0 = V_0^+ \quad V_0 = \Pi(V_0)$$

$$H + \varepsilon \hat{V}(\varepsilon) = (1 - i \varepsilon K_1) (H + \varepsilon V) (1 + i \varepsilon K_1) + O(\varepsilon^2)$$

$$V_0 = i [H, K_1] + V$$

$$\Pi([H, K_1]) = \sum_k P_k (HK_1 - K_1 H) P_k = \sum_k P_k (e_k K_1 - K_1 e_k) P_k = 0$$

$$\Rightarrow V_0 = \Pi(V_0) = \Pi(V) = V_Z \quad 1) \quad V_0 = \Pi(V) = V_Z$$

$$i[H, K_1] = -(V - \Pi(V)) = -V_{\text{off}} \quad 2) \quad i[H, K_1] = -V_{\text{off}}$$

$$V_{\text{off}} = V - \Pi(V) = \sum_{k, l: k \neq l} P_k V P_l = \sum_k P_k V (1 - P_k) \quad \begin{matrix} \text{off-diagonal} \\ \text{part of } V \end{matrix}$$

Homological equation for K_1

$$i[H, K_1] = V_{\text{off}} \quad (*)$$

$$i P_k (H K_1 - K_1 H) P_j = -P_k \vee P_j \quad k \neq j$$

$$i (\ell_k - \ell_j) P_k K_1 P_j = -P_k \vee P_j$$

$$P_k H P_j = i \frac{1}{\ell_k - \ell_j} P_k \vee P_j$$

$$K_1 = i \sum_{k \neq j} \frac{1}{\ell_k - \ell_j} P_k \vee P_j + \pi(K_1) = i \sum_j S_j \vee P_j + \pi(K_1)$$

$$S_j = \sum_{k: k \neq j} \frac{1}{\ell_k - \ell_j} P_k \quad \text{Reduced resolvent at } \ell_j : \\ (H - \ell_j) S_j = S_j (H - \ell_j) = 1 - P_j$$

Equation $(*)$ has a unique solution with $\pi(K_1) = 0$:

$$K_1 = i \sum_{j=1}^m S_j \vee P_j$$

$$U_\varepsilon = 1 + i \varepsilon K_1 + O(\varepsilon^2)$$

$$\mathcal{S}_\infty = \sup_{t \in \mathbb{R}} \| e^{it(H+\varepsilon V)} - e^{it(H+\varepsilon \hat{V})} \| \leq 2 \| U_\varepsilon - 1 \| = 2 \varepsilon \| K_1 \| + O(\varepsilon^2)$$

$$\|K_1\|^2 = \|K_1 K_1^*\| = \left\| \sum_j S_j V P_j V S_j \right\| \leq \sum_j \|S_j V P_j V S_j\| \leq \sum_j \|S_j\|^2 \|V\|^2$$

$$\|S_j\| = \left\| \sum_{k: k \neq j} \frac{1}{\ell_k - \ell_j} P_k \right\| = \max_{k: k \neq j} \left| \frac{1}{\ell_k - \ell_j} \right| \leq \frac{1}{\eta}$$

$$\gamma = \min_{k, j: k \neq j} |\ell_k - \ell_j| \quad \text{minimal spectral gap}$$

$$\|K_1\|^2 \leq \sum_{j=1}^m \frac{1}{\eta^2} \|V\|^2 = \frac{m}{\eta^2} \|V\|^2$$

$$\mathcal{S}_\infty \leq 2 \frac{\varepsilon \sqrt{m}}{\eta} \|V\| + O(\varepsilon^2) \leq \frac{7 \sqrt{m}}{\eta} \varepsilon \|V\|$$

Conclusions

Definition

- $S \in \mathcal{H}^1$ fragile if $\exists V = V^+ \in B(\mathcal{H})$ s.t.

$$\sup_{t \in \mathbb{R}} \| e^{it(H+\varepsilon V)} S e^{-it(H+\varepsilon V)} - S \| = O(1)$$

$$\left[\inf_{\varepsilon > 0} \sup_{t \in \mathbb{R}} \| e^{it(H+\varepsilon V)} S e^{-it(H+\varepsilon V)} - S \| > 0 \right]$$

- $S \in \mathcal{H}^1$ robust if it is not fragile.

$$\left[\forall V = V^+ \in B(\mathcal{H}) : \sup_{t \in \mathbb{R}} \| e^{it(H+\varepsilon V)} S e^{-it(H+\varepsilon V)} - S \| = O(\varepsilon) \right]$$

Theorem [Robust symmetries]

$S \in \{H\}^{\dagger}$ symmetry of $H = H^+$. Then the following assertions are equivalent:

1. S is robust $\left[\sup_{t \in \mathbb{R}} \|e^{it(H+\varepsilon V)} S e^{-it(H+\varepsilon V)} - S\| = O(\varepsilon) \right.$
 $\forall V = V^+, \|V\| = 1 \left. \right]$

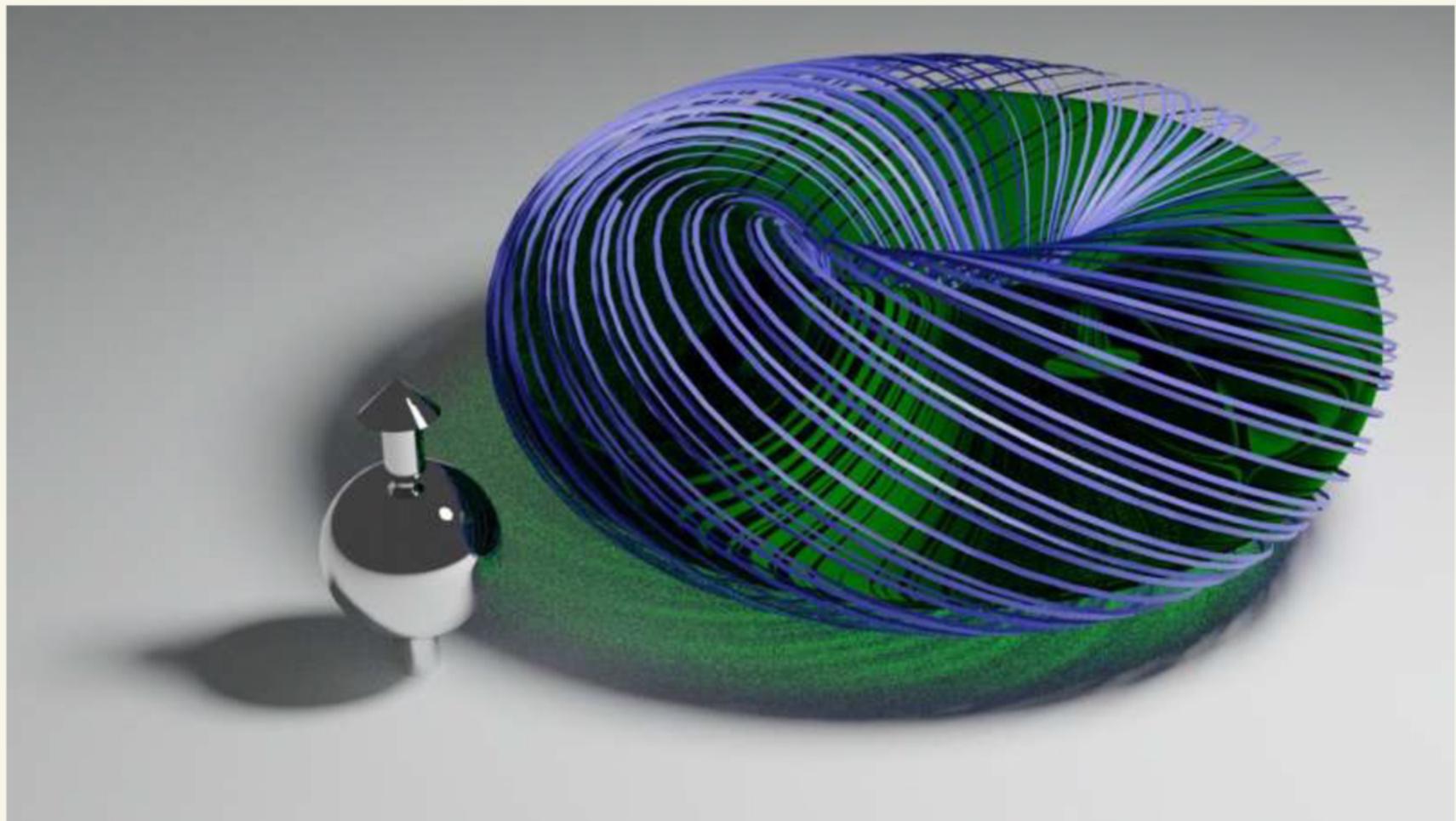
2. $S \in \{H\}^{\dagger\dagger}$

3. $\exists C = C_H > 0$ s.t. $\forall V = V^+ \in B(\mathcal{H}) :$

$$\sup_{t \in \mathbb{R}} \|e^{it(H+\varepsilon V)} S e^{-it(H+\varepsilon V)} - S\| \leq C_H \varepsilon \|V\| \|S\|$$

$$\left[C_H < 14 \frac{\sqrt{m}}{\eta} , \quad H = \sum_{k=1}^m e_k P_k , \quad \eta = \min_{k \neq j} |e_k - e_j| \right]$$

Thank you!



References

- D. Burgarth, P.F., H. Nakazato, S. Pascazio, R. Yuasa
Generalized adiabatic theorem and strong coupling limits
Quantum 3, 152 (2019)
- D. Burgarth, P.F., H. Nakazato, S. Pascazio, R. Yuasa
Eternal adiabaticity in quantum evolution
Phys. Rev. A 103, 032214 (2021)
- D. Burgarth, P.F., H. Nakazato, S. Pascazio, R. Yuasa
*Kolmogorov-Arnold-Moser stability for conserved quantities
in finite-dimensional quantum systems*
Phys. Rev. Lett. 126, 150401 (2021)