Fast is hot: energetics of information erasure and the overhead to Landauer's bound at low dissipation



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Energetic cost of information processing

- Our brain: ~20 W (vs ~100 W for our body)
- Internet: 10% of electricity worldwide
- Supercomputers: costs are driven by Flops/W



Energetic cost of information processing

Physical Limit ? Rolf Landauer (IBM, 1961):

1-bit erasure costs at least $k_B T \ln 2 = 3 \times 10^{-21} \text{ J}$



R. Landauer, *IBM Journal of Research and Development* **5**, 183 (1961)

Energetic cost of information processing

T

Physical Limit ? **Rolf Landauer** (IBM, 1961): 1-bit erasure costs at least $k_B T \ln 2 = 3 \times 10^{-21} \text{ J}$

$$S = k_B \ln N$$
 [Boltzmann]

$$N = 2 \qquad \qquad N = 1$$
$$\Delta S = -k_B \ln 2 \ge -\frac{\langle \mathcal{Q} \rangle}{T}$$

 $\{0,1\}$ erasure $\{0\}$

W memory Q

$$\langle \mathcal{Q} \rangle = \langle \mathcal{W} \rangle \ge k_B T \ln 2$$



R. Landauer, *IBM Journal of Research and Development* **5**, 183 (1961)

Roadmap

Physical Limit ? **Rolf Landauer** (IBM, 1961): 1-bit erasure costs at least $k_B T \ln 2 = 3 \times 10^{-21} \text{ J}$

> Experimental demonstration ? Distance to the bound ?

Fast operation ?

Energetics ?

$$\langle \mathcal{Q} \rangle = \langle \mathcal{W} \rangle \gtrsim k_B T \ln 2$$







L. Bellon - University of Helsinki - may 2021





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- 1 DOF
- Bistable potential
- k_BT scale
- Tunable potential $\longrightarrow U(x,\lambda)$
- \bullet Measure $\mathcal{Q}\,\&\,\mathcal{W}$

/ white noise, variance 1

Overdamped Langevin eq.

$$\dot{x} + \frac{\partial U}{\partial x} = \sqrt{2k_B T \gamma} \eta$$

$$\mathcal{W} = \int_0^\tau \frac{\partial U}{\partial \lambda} \dot{\lambda} dt \qquad \mathcal{Q} = -\int_0^\tau \frac{\partial U}{\partial x} \dot{x} dt$$

$$\Delta U = \mathcal{W} - \mathcal{Q}$$

K. Sekimoto, Stochastic Energetics, Lecture Notes in Physics 799 (Springer, 2010)

Stochastic thermodynamics



white noise, variance 1

Underdamped Langevin eq. $m\ddot{x} + \gamma\dot{x} + \frac{\partial U}{\partial x} = \sqrt{2k_BT\gamma}\eta$

$$\mathcal{W} = \int_0^\tau \frac{\partial U}{\partial \lambda} \dot{\lambda} dt \qquad \mathcal{Q} = -\int_0^\tau \frac{\partial U}{\partial x} \dot{x} dt - \Delta K$$

$$\Delta K + \Delta U = \mathcal{W} - \mathcal{Q}$$

S. Albert, A. Archambault, A. Petrosyan, C. Crauste-Thibierge, L. Bellon, S. Ciliberto, EPL 131, 10008 (2020)

• 1 DOF

- Bistable potential
- k_BT scale
- Tunable potential
- ✓ Measure Q & W





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Karel Proesmans, Jannik Ehrich, John Bechhoefer, Phys. Rev. Lett. 125, 100602 (2020)



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Underdamped oscillator

- 1 DOF
- Bistable potential
- k_BT scale
- Tunable potential
- \bullet Measure $\mathcal{Q}\,\&\,\mathcal{W}$



$$U(x,\lambda) = \frac{1}{2}k(x-\lambda)^{2}$$

$$\sigma = \sqrt{\frac{k_{B}T}{k}} \sim 1 \text{ nm}$$

$$f_{0} = \sqrt{\frac{k}{m}} = 1270 \text{ Hz}, \quad Q \sim 10 \rightarrow \tau_{\text{relax}} = \frac{Q}{\pi f_{0}} \sim 2.5 \text{ms}$$

S. Albert, A. Archambault, A. Petrosyan, C. Crauste-Thibierge, L. Bellon, S. Ciliberto, EPL 131, 10008 (2020)

Underdamped oscillator

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- \bullet Measure $\mathcal{Q}\,\&\,\mathcal{W}$

 $U(x,\lambda)$ arbitrary !

$$\sigma = \sqrt{\frac{k_B T}{k}} \sim 1 \,\mathrm{nm}$$

$$Q$$

 $f_0 = \sqrt{\frac{\pi}{m}} = 1270 \,\mathrm{Hz}, \quad Q \sim 10 \rightarrow \tau_{\mathrm{relax}} = \frac{\tau}{\pi f_0} \sim 2.5 ms$

Bistable underdamped oscillator



<u>Click here to play video online</u>





$$\langle \mathcal{Q} \rangle = \langle \mathcal{W} \rangle = k_B T \left(\ln 2 + \frac{B}{\tau} \right)$$

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S. Dago, J. Pereda, N. Barros, S. Ciliberto, L. Bellon, Phys. Rev. Lett. 126, 170601 (2021)

$$P(x,v) = \frac{1}{Z}e^{-\frac{1}{2}\beta mv^2}e^{-\frac{1}{2}\beta k(|x|-\lambda)^2}, \quad \beta = \frac{1}{k_B T}, \quad Z = \frac{2\pi}{\sqrt{km\beta}}V, \quad V = 1 + \operatorname{erf}\left(\sqrt{\frac{k\beta}{2}}\lambda\right)$$

$$\langle E \rangle = \langle U + K \rangle = -\frac{\partial \ln Z}{\partial \beta} = k_B T + k_B T^2 \frac{\partial \ln V}{\partial T}$$

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$$\frac{\partial \langle E \rangle}{\partial T} \dot{T} + \frac{\partial \langle E \rangle}{\partial \lambda} \dot{\lambda} = \langle \dot{E} \rangle = \langle \dot{\mathcal{W}} \rangle - \langle \dot{\mathcal{Q}} \rangle = -k_B T \frac{\partial \ln V}{\partial \lambda} \dot{\lambda} + \frac{\omega_0}{Q} k_B (T - T_0)$$

S. Dago, L. Bellon, arXiv: 2105.12023 (2021)

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S. Dago, L. Bellon, arXiv: 2105.12023 (2021)

<u>Click here to play</u> <u>video online</u>

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Adiabatic limit

$$Q \gg 1 \longrightarrow \langle \dot{\mathcal{Q}} \rangle = -\frac{\omega_0}{Q} k_B (T - T_0) \sim 0$$

$$\left[k_B \ln Z + \frac{\langle E \rangle}{T}\right]_0^\tau = \Delta S = \int_0^\tau \frac{\langle \dot{Q} \rangle}{T} dt = 0$$

$$\langle E \rangle = k_B T + k_B T^2 \frac{\partial \ln V}{\partial T}$$

$$Z = \frac{2\pi}{\sqrt{km}} k_B T V$$

$$V = 1 + \operatorname{erf} \left(\sqrt{\frac{k\beta}{2}} \lambda \right)$$

$$TV = T_0 V_0$$

$$T(\tau) = 2T_0$$

$$\langle \mathcal{W} \rangle = \Delta \langle E \rangle = k_B T_0$$

Fast is hot, but not too hot !

$$Q \gg 1 \longrightarrow \langle \mathcal{W} \rangle = \langle \mathcal{Q} \rangle \sim k_B T_0$$
(delayed)

Underdamped stochastic thermodynamics

- Fast operation, high statistics
- Illustration of Landauer's bound
- $k_B T_0 \ln 2$ comes from compression
- Fast is hot
- Full energetics description
- Adiabatic limit: $\langle \mathcal{W} \rangle = \langle \mathcal{Q} \rangle = k_B T_0$

Thank you for your attention

Kiitos huomiostasi :-)

[according to Google translate]

To go further...

R. Landauer, *IBM Journal of Research and Development* **5**, 183 (1961) doi: <u>10.1147/rd.53.0183</u>

A. Bérut, A. Arakelyan, A. Petrosyan, S. Ciliberto, E. Dillenschneider, E. Lutz, *Nature* **483**, 187 (2012) doi: <u>10.1038/nature10872</u>

S. Dago, J. Pereda, N. Barros, S. Ciliberto, L. Bellon, *Phys. Rev. Lett.* **126**, 170601 (2021) doi: <u>10.1103/PhysRevLett.126.170601</u>

S. Dago, L. Bellon (2021) arXiv: <u>2105.12023</u> [cond-mat.stat-mech]

http://perso.ens-lyon.fr/ludovic.bellon

