# Inverse problems for wave equation

# Matti Lassas



HELSINKI UNIVERSITY OF TECHNOLOGY Institute of Mathematics



Finnish Centre of Excellence in Inverse Problems Research Motivation Let  $\Omega \subset \mathbb{R}^m$ ,

u(x,t) satisfy a wave equation in  $\Omega \times \mathbb{R}$ 

#### **Inverse problem:**

Can we determine the coefficients of the wave equation, i.e., physical model in  $\Omega$  by observing

u(x,t) near  $\partial\Omega \times \mathbb{R}$ 

for all possible solutions u(x, t)?

The inverse problem has no unique solution as

We can change definition of x-coordinate: Let

$$v(x,t) = u(\phi(x),t)$$

where

$$\phi: \Omega \to \Omega, \quad \phi|_{\partial\Omega} = id$$

We can change scale of *u*-coordinate: Let

$$w(x,t) = \kappa(x)u(x,t)$$

where  $\kappa(x) > 0$ .

All functions u, v and w model the same physical process.

Let us consider  $\Omega$  as Riemannian manifold

 $d_g(x,y) =$  travel time between x and y.

Let us identify all isometric Riemannian manifolds, that is, we ask following question

Do the boundary measurements determine uniquely the isometry type of the Riemannian manifold?

Let *u* satisfy the wave equation

 $u_{tt} + a(x, D)u = 0.$ 

Then the gauge transformation of u,

$$w(x,t) = \kappa(x)u(x,t)$$

satisfy

$$w_{tt} + a_{\kappa}(x, D)w = 0,$$

where

$$a_{\kappa}(x,D)w = \kappa a(x,D)(\kappa^{-1}w)$$

We say that the gauge equivalence class of a(x, D) is

$$[a(x,D)] = \{a_{\kappa}(x,D) : \kappa > 0\}$$

Can the equivalence class be uniquely determined?

# 1 Setting of the problem I

Let us consider the wave equation

$$\begin{aligned} u_{tt}(x,t) + Au(x,t) &= 0, & \text{in } M \times \mathbb{R}_+, \\ u_{t=0} &= 0, & u_t|_{t=0} = 0, \\ u_{\partial M \times \mathbb{R}^+} &= f \end{aligned}$$

where M is a m-dimensional manifold and

$$Au = -\sum_{j,k=1}^{m} a^{jk} \frac{\partial^2 u}{\partial x^i \partial x^j} + \sum_{j=1}^{m} b^j \frac{\partial u}{\partial x^j} + cu,$$

where  $a^{jk}, b^j, c$  are real, smooth,  $[a^{jk}(x)] > 0$ . In addition ... Assume that there is dV such that A is selfadjoint in  $L^2(M, dV)$  with

$$\mathcal{D}(A) = H^2(M) \cap H^1_0(M).$$

Now

$$g^{jk} = a^{jk}$$
 defines a metric tensor on  $M$ .

This makes (M, g) a Riemannian manifold.

# **1.1** Invariant inverse problem

The Robin-to-Dirichlet map is

$$\Lambda : (\partial_{\nu} u + \sigma u)|_{\partial M \times \mathbb{R}_+} \mapsto u|_{\partial M \times \mathbb{R}_+}.$$

## **Dynamical inverse problem:**

Let  $\partial M$  and the map  $\Lambda$  be given. Can we determine

(M, g) and [A(x, D)]?

**Energy flux through boundary** The energy of the wave at time *t* is

$$E(u,t) = \int_{M} \left( |\partial_t u(t)|^2 + |\operatorname{Grad} u(t)|_g^2 + q|u(t)|^2 \right) dV + \int_{\partial M} \sigma |u(t)|^2 dS.$$

For  $h = u|_{\partial M \times \mathbb{R}_+} \in C_0^\infty(\partial M \times \mathbb{R}_+)$  let

$$\Pi(h) = \lim_{t \to \infty} E(u, t).$$

### **Inverse problem for energy flux:**

Let  $\partial M$  and map  $\Pi$  be given. Can we determine

(M,g) and [A(x,D)]?

#### **Inverse boundary spectral problem:**

Operator A has in  $L^2(M, dV)$  orthonormal eigenfunctions  $\varphi_j$ ,

$$(-\Delta_g + P + q - \lambda_j)\varphi_j = 0,$$
  
$$\partial_{\nu}\varphi_j|_{\partial M} = 0.$$

Let boundary spectral data

$$\{\partial M, \lambda_j, \varphi_j|_{\partial M}, j = 1, 2, \dots\}$$

be given. Can we determine

(M,g) and [A(x,D)]?

- The above inverse problems are equivalent.
- Consider gauge equivalence class [A(x, D)] of operator A(x, D). Then there is a unique Schrödinger operator

$$-\Delta_g + q \in [A(x,D)].$$

Because of this we next restrict ourselves to the case  $A = -\Delta_g + q$ .

# 2 Setting of the problem II

Denote by

$$u^f = u^f(x,t)$$

the solutions of

$$\begin{aligned} u_{tt} - \Delta_g u + qu &= 0 \quad \text{on } M \times \mathbb{R}_+, \\ -\partial_\nu u|_{\partial M \times \mathbb{R}_+} &= f, \\ u|_{t=0} &= 0, \quad u_t|_{t=0} = 0, \end{aligned}$$

where  $\nu$  is unit interior normal of  $\partial M$ . Define

$$\Lambda_T f = u^f |_{\partial M \times (0,T)}.$$

We denote  $\Lambda = \Lambda_{\infty}$ . Assume that we are given the boundary data  $(\partial M, \Lambda)$ .

Results on the problem:

- Nachman-Sylvester-Uhlmann '88.
- $c(x)^2 \Delta$  in  $\mathbb{R}^m$  by boundary control method, Belishev '87, Belishev-Kurylev '87.
- $\Delta_g$  on manifold, Belishev-Kurylev '92.
- Local controllability, Tataru '95.
- Equivalence of above inverse problems Katchalov-Kurylev-L.-Mandache 2004
- Maxwell's equations Kurylev-L.-Somersalo 2006.
- Dirac system Kurylev-L.-Somersalo 2006.
- Reconstruction based on iterated time reversal Bingham-Kurylev-L.-Siltanen 2007.

In the following we consider the geometric version of the Belishev-Kurylev-Tataru method, or Boundary Control method, see references [1-7].

## 2.1 Blagovestchenskii identity

## Lemma 2.2 Let $f, h \in L^2(\partial M \times [0, 2T])$ . Then

$$\int_M u^f(x,T)u^h(x,T)\,dV_\mu(x) =$$

 $\int_{[0,2T]^2} \int_{\partial M} J(t,s) \left[ f(t)(\Lambda_{2T}h)(s) - (\Lambda_{2T}f)(t)h(s) \right] dS_g(x) dt ds,$ 

where  $J(t,s) = \frac{1}{2}\chi_L(s,t)$  and  $\chi_L$  being the characteristic function of the triangle

 $L = \{ (s, t) \in \mathbb{R}_+ \times \mathbb{R}_+ : t + s \le 2T, s < t \}.$ 

**Proof.** Let  $w(t,s) = \int_M u^f(t) u^h(s) dV_{\mu}$ . Integrating by parts, we see that

$$\begin{aligned} (\partial_t^2 - \partial_s^2)w(t,s) &= -\int_M \left[Au^f(t)u^h(s) - u^f(t)Au^h(s)\right] dV_\mu(x) \\ &= -\int_{\partial M} \left[\partial_\nu u^f(t)u^h(s) - u^f(t)\partial_\nu u^h(s)\right] dS_g \\ &= \int_{\partial M} \left[f(t)\Lambda h(s) - \Lambda f(t)h(s)\right] dS_g. \end{aligned}$$

Moreover,

$$w|_{t=0} = w|_{s=0} = 0, \quad \partial_t w|_{t=0} = \partial_s w|_{s=0} = 0.$$

Thus we can find w(s,t) by solving a wave equation with known initial data and right side.

## 2.3 Domains of influence

Let  $\Gamma \subset \partial M$  be a non-empty open set. We denote by  $L^2(\Gamma \times [0,T])$  the subspace of  $L^2(\partial M \times [0,T])$  that consists of the functions f with supp  $(f) \subset \overline{\Gamma} \times [0,T]$ . Definition 2.4 The subset  $M(\Gamma, \tau) \subset M, \tau > 0$ ,

 $M(\Gamma, \tau) = \{ x \in M : d(x, \Gamma) \le \tau \}$ 

is called the domain of influence of  $\Gamma$  at time  $\tau$ .



Lemma 2.5 Let  $f \in L^2(\Gamma \times [0,T])$ . Then

supp  $(u^f(\tau)) \subset M(\Gamma, \tau).$ 

**Proof.** The result follows finite velocity of wave propagation.



We denote by  $L^2(\Omega)$ ,  $\Omega \subset M$ , the subspace of  $L^2(M)$ , which consists of all functions  $f \in L^2(M)$  that are equal to zero in  $M \setminus \Omega$ . We prove following Tataru-type controllability type theorem.

**Theorem 1** Let  $\tau > 0$ . The linear subspace,

$$\{u^f(\tau) \in L^2(M(\Gamma,\tau)): f \in C_0^\infty(\Gamma \times [0,\tau])\},\$$

is dense in  $L^2(M(\Gamma, \tau))$ .

**Proof.** Let  $\psi \in L^2(M(\Gamma, \tau))$  be such that

$$\langle u^f(\cdot,\tau),\psi\rangle=0$$

### for all $f \in C_0^{\infty}(\Gamma \times [0, \tau])$ . To prove the claim, it is sufficient to show that $\psi = 0$ .

We consider the wave equation,

$$(\partial_t^2 - \Delta_g + q)e = 0, \quad \text{in} \quad M \times (0, \tau),$$
  
$$\partial_{\nu} e|_{\partial M \times (0, \tau)} = 0, \quad e|_{t=\tau} = 0, \quad \partial_t e|_{t=\tau} = \psi.$$

Integrating by parts we obtain

$$0 = \int_{M \times (0,\tau)} [u^f (\partial_t^2 - \Delta_g + q)e - ((\partial_t^2 - \Delta_g + q)u^f)e] dV_g dt$$
  
$$= \int_M u^f(\tau) \psi dV_g + \int_{\partial M \times (0,\tau)} f e dS_g dt$$
  
$$= \int_{\partial M \times (0,\tau)} f e dS_g dt,$$

for all  $f \in C_0^{\infty}(\Gamma \times [0, \tau])$ . This yields that the Cauchy data of *e* vanish on  $\Gamma \times (0, \tau)$ . Recall that  $e(x, \tau) = 0$ . We continue e onto  $t \in [\tau, 2\tau]$  as

$$E(x,t) = \begin{cases} e(x,t), & \text{for } t \leq \tau, \\ -e(x,2\tau-t), & \text{for } t > \tau. \end{cases}$$

Then  $E \in C([0, 2\tau]; H^1(M)) \cap C^1([0, 2\tau]; L^2(M))$  and

$$(\partial_t^2 - \Delta_g + q)E = 0$$
 in  $M \times (0, \tau)$ .

The Cauchy data of E vanish on  $\Gamma \times ([0, 2\tau] \setminus \{\tau\})$ . Since  $\partial_{\nu} E \in L^2(\partial M \times (0, 2\tau))$ , we see that

$$E|_{\Gamma \times (0,2\tau)} = 0, \quad \partial_{\nu} E|_{\Gamma \times (0,2\tau)} = 0.$$

Then  $\psi = 0$  by the following Tataru-Holmgren-John theorem.

**Theorem 2** Let u be a solution in  $M \times (0, 2\tau)$  of the wave equation

$$(\partial_t^2 - \Delta_g + q)u = 0$$
 in  $M \times (0, 2\tau)$ .

such that for an open set  $\Gamma \subset \partial M$ ,

$$u|_{\Gamma \times [0,2\tau]} = 0, \ \partial_{\nu} u|_{\Gamma \times (0,2\tau)} = 0.$$

Then, at  $t = \tau$ , the function u and its derivative  $\partial_t u$  vanish in the domain of influence of  $\Gamma$ ,

$$u(x,\tau) = 0, \ \partial_t u(x,\tau) = 0 \quad \text{ for } x \in M(\Gamma,\tau).$$



## 2.6 Wave basis

The set

$$\{u^f(\tau) \in L^2(M(\Gamma,\tau)): f \in L^2(\Gamma \times [0,\tau])\}$$

is dense in  $L^2(M(\Gamma, \tau))$ . Thus, there are functions  $f_j$ ,  $j = 1, 2, \ldots$ , such that  $\{u^{f_j}(\tau)\}_{j=1}^{\infty}$  form an orthonormal basis in the space  $L^2(M(\Gamma, \tau))$ . We will construct such functions  $f_j$  from the boundary data. The corresponding basis  $\{u^{f_j}(\tau)\}_{j=1}^{\infty}$  is called the wave basis. **Lemma 2.7** Let  $\tau > 0$ . Given the boundary data it is possible to construct boundary sources  $f_j \in L^2(\Gamma \times [0, \tau])$  such that

$$v_j = u^{f_j}(\tau), \ j = 1, 2, \dots,$$

form an orthonormal basis of  $L^2(M(\Gamma, \tau))$ .



**Proof.** Let  $\{h_j\}_{j=1}^{\infty} \subset C_0^{\infty}(\Gamma \times (0, \tau))$  be a complete set in  $L^2(\Gamma \times [0, \tau])$ . We can compute that inner products

$$c_{jk} = \langle u^{h_j}(\tau), u^{h_k}(\tau) \rangle.$$

Next we use the Gram-Schmidt orthogonalization procedure to construct  $f_j$ . More precisely, we define  $f_j \in L^2(\Gamma \times [0, \tau])$  recursively by

$$g_j = h_j - \sum_{k=1}^{j-1} \langle u^{h_j}(\tau), u^{f_k}(\tau) \rangle f_k,$$

$$f_j = \frac{g_j}{\langle u^{g_j}(\tau), u^{g_j}(\tau) \rangle^{1/2}}.$$

When  $g_j = 0$ , we remove the corresponding  $h_j$  from the original sequence and continue the procedure.

Since  $\{h_j\} \subset C_0^{\infty}(\Gamma \times [0, \tau])$ , we have  $f_j \in C_0^{\infty}(\Gamma \times [0, \tau])$ . Thus  $u^{f_j}(\tau) \in C^{\infty}(M)$ .

Let T > diam(M). Then  $M(\partial M, T) = M$ , and the corresponding wave basis

$$\{u^{\eta_j}(T)\}_{j=1}^\infty$$

is the orthonormal basis in  $L^2(M)$ . Next we reserve the notation  $\eta_j \in C^{\infty}(\partial M \times (0,T))$  for such boundary values.

## 2.8 **Projectors**

Denote by  $P_{\Gamma,\tau}$  the orthogonal projector in  $L^2(M)$  onto the space  $L^2(M(\Gamma,\tau))$ ,

$$P_{\Gamma,\tau}: L^2(M) \to L^2(M(\Gamma,\tau)),$$
$$(P_{\Gamma,\tau}a)(x) = \chi_{M(\Gamma,\tau)}(x)a(x),$$

where  $\chi_{M(\Gamma,\tau)}$  is the characteristic function of the domain of influence  $M(\Gamma,\tau)$ ,

$$\chi_{M(\Gamma,\tau)}(x) = \begin{cases} 1, & \text{for } x \in M(\Gamma,\tau), \\ 0, & \text{for } x \notin M(\Gamma,\tau). \end{cases}$$

**Lemma 2.9** Let  $f, h \in L^2(\partial M \times [0,T])$  and  $\Gamma \subset \partial M$  be an open set. Then, given the the map  $\Lambda$ , it is possible to find the inner product

$$\langle P_{\Gamma,\tau} u^f(t), u^h(s) \rangle = \int_{M(\Gamma,\tau)} u^f(x,t) u^h(x,s) dV_g$$

for any  $0 \le t, s, \tau \le T$ .



**Proof.** We can find  $f_j \in C_0^{\infty}(\Gamma \times [0, \tau])$  such that  $v_j = u^{f_j}(\tau)$  is an orthonormal basis in  $L^2(M(\Gamma, \tau))$ , Then, for any  $a \in L^2(M(\Gamma, \tau))$ ,

$$a = \sum_{j=1}^{\infty} \langle a, v_j \rangle \, v_j.$$

As  $\langle P_{\Gamma,\tau}u^f(t), v_j \rangle = \langle u^f(t), v_j \rangle$ , we have

$$\langle P_{\Gamma,\tau} u^f(t), u^h(s) \rangle = \sum_{j=1}^{\infty} \langle u^f(t), v_j \rangle \langle u^h(s), v_j \rangle.$$

Here  $\langle u^f(t), v_j \rangle$  and  $\langle u^h(s), v_j \rangle$  can be computed using boundary data.

Denote by  $M(y,\tau)$  the domain of influence of a point  $y \in \partial M$ ,

$$M(y,\tau) = \{ x \in M : \ d(x,y) \le \tau \},\$$

and by  $P_{y,\tau}$  the orthoprojector

$$P_{y,\tau}: L^2(M) \to L^2(M(y,\tau)).$$

**Corollary 2.10** Let  $f, h \in L^2(\partial M \times [0,T])$  and  $y \in \partial M$  be given. Then the boundary data determine the inner product

$$\langle P_{y,\tau} u^f(t), u^h(s) \rangle = \int_{M(y,\tau)} u^f(x,t) u^h(x,s) dV_g$$

for any  $0 \le t, s, \tau \le T$ .

**Proof.** Let  $\Gamma_l$ , l = 1, 2, ... be open sets such that

$$\Gamma_{l+1} \subset \Gamma_l, \quad \bigcap_{l=1}^{\infty} \Gamma_l = \{y\}.$$

Then,

$$\lim_{l \to \infty} \chi_{M(\Gamma_l, \tau)}(x) = \chi_{M(y, \tau)}(x)$$

pointwise. By the Lebesgue dominated convergence theorem,

$$\lim_{l \to \infty} \langle P_{\Gamma_l,\tau} u^f(t), u^h(s) \rangle = \langle P_{y,\tau} u^f(t), u^h(s) \rangle.$$

**Corollary 2.11** Let  $f \in L^2(\partial M \times [0,T])$  and  $y \in \partial M$ . Then the boundary data determine uniquely the inner product

$$\langle P_{y,\tau}u^{\eta_k}(T), u^{\eta_l}(T) \rangle = \sum_{j=1}^{\infty} \langle u^{\eta_k}(T), u^{f_j}(\tau) \rangle \langle u^{\eta_l}(T), u^{f_j}(t) \rangle,$$

where  $\{u^{f_j}(\tau)\}_{j=1}^{\infty}$  form an orthonormal basis in  $L^2(M(y,\tau))$ .

**Corollary 2.12** Let  $f \in L^2(\partial M \times [0,T])$  and  $y_j \in \partial M$ ,  $\tau_j > 0$ . Then the boundary data determine the inner product

$$\langle Q_N u^f(s), u^{\eta_l}(T) \rangle$$

where

$$Q_N = \prod_{j=1}^N P_{y_j,\tau_j}$$

and  $\{u^{f_j}(\tau)\}_{j=1}^{\infty}$  form an orthonormal basis in  $L^2(M(y,\tau))$ .

**Proof.** For N = 1 the claim follows from Corollary 2.11. Assume now that it is valid for N - 1. We can write

$$Q_{N-1}u^f(s) = \sum_{k=1}^{\infty} \langle Q_{N-1}u^f(s), u^{\eta_k}(T) \rangle u^{\eta_k}(T)$$

and

$$\langle Q_N u^f(T), u^{\eta_l}(T) \rangle = \langle P_{y_N, \tau_N} Q_{N-1} u^f(T), u^{\eta_l}(T) \rangle$$
$$= \sum_{k=1}^{\infty} \langle P_{y_N, \tau_N} u^{\eta_k}(T), u^{\eta_l}(T) \rangle \langle Q_{N-1} u^f(s), u^{\eta_k}(T) \rangle.$$

From this the claim follows by induction.
Observations:

• We can compute the Gram matrix  $[q_{jk}]_{j,k=1}^{\infty}$ ,

$$q_{jk} = \langle Qu^{\eta_j}(T), u^{\eta_k}(T) \rangle$$

where  $\{u^{\eta_j}(T)\}_{j=1}^{\infty}$  is an orthonormal basis in  $L^2(M)$  and

$$Q = \left(\prod_{j=1}^{N} P_{y_j,\tau_j^+}\right) \left(\prod_{j=1}^{N} (1 - P_{y_j,\tau_j^-})\right)$$

• The projector  $Q: L^2(M) \to L^2(M)$  is

$$Qv(x) = \chi_I(x) v(x), \quad I = \bigcap_{j=1}^N (M(y_j, \tau_j^+) \setminus M(y_j, \tau_j^-)).$$

• The projector  $Q: L^2(M) \to L^2(M)$  is

$$Qv(x) = \chi_I(x) v(x), \quad I = \bigcap_{j=1}^N (M(y_j, \tau_j^+) \setminus M(y_j, \tau_j^-)).$$

■ The projector  $Q: L^2(M) \to L^2(M)$  vanishes, that is, its Gram matrix is zero if and only if

$$m(I) = 0, \quad I = \bigcap_{j=1}^{N} (M(y_j, \tau_j^+) \setminus M(y_j, \tau_j^-)).$$

Thus we can check using boundary data if m(I) = 0.



Boundary distance functions. For  $x \in M$  define  $r_x(y) = d(x, y), y \in \partial M$ . Let

$$R: M \to C(\partial M), \quad R(x) = r_x.$$

Next we consider R(M) as a submanifold on  $C(\partial M)$ . **Theorem 3** Using boundary data we can determine

$$R(M) = \{ r_x \in C(\partial M) : x \in M \}.$$

Thus the constructed set R(M) can be identified with M.

By previous observations, it is enough to prove the following result:

**Lemma 2.13** Let  $\{z_n\}_{n=1}^{\infty}$  be a dense set on  $\partial M$ . Then  $r(\cdot) \in C(\partial M)$  lies in R(M) if and only if, for any N > 0,

$$I_N = \bigcap_{n=1}^N M(z_n, r(z_n) + \frac{1}{N}) \cap \bigcap_{n=1}^N (M(z_n, r(z_n) - \frac{1}{N}))^c.$$

satisfies

$$m(I_N) \neq 0 \tag{1}$$

Moreover, condition (1) can be verified using the boundary data.

**Proof** *"If"-part.* Assume that  $r(\cdot) = r_x(\cdot)$  with some  $x \in M$ . Consider a ball  $B_{1/N}(x)$ . Then,

$$B_{1/N}(x) \subset M(z, r(z) + \frac{1}{N}) \setminus M(z, r(z) - \frac{1}{N}).$$

Thus if  $B_{1/N}(x) \subset I_N$  and  $m(I_N) \neq 0$ .

"Only if"-part. Assume that  $m(I_N) \neq 0$ . Then there exists

$$x_N \in \bigcap_{n=1}^N \left( M(z_n, r(z_n) + \frac{1}{N}) \setminus M(z_n, r(z_n) - \frac{1}{N}) \right).$$

Since *M* is compact, we can choose a subsequence of  $x_N$  (denoted also by  $x_N$ ), so that there exists a limit

$$x = \lim_{n \to \infty} x_N.$$

By continuity of the distance function, it follows from (2) that

$$d(x, z_n) = r(z_n), \quad n = 1, 2, \dots$$

Since  $\{z_n\}$  are dense in  $\partial M$ , we see that r(z) = d(x, z) for all  $z \in \partial M$ . Thus  $r = r_x$ .

















# **2.14** Reconstruction of (M, g) from R(M).

**Theorem 4** The set R(M) has a Riemannian manifold structure which is isometric to (M, g).

Recall that for  $x \in M$ 

$$r_x(z) = d(x, z), \ z \in \partial M$$

and that

$$R: M \to C(\partial M), \quad R(x) = r_x.$$

Next we consider R(M) as a submanifold on  $C(\partial M)$ .



By triangular inequality we have

$$||r_x - r_y||_{C(\partial M)} \le d(x, y), \quad x, y \in M.$$

**Example:** Consider that case when all geodesics of a compact manifold (M, g) are the shortest curves between their endpoints and all geodesics can be continued to geodesics that hit the boundary. Then for any  $x, y \in M$  the geodesic from x to y hits later to  $z \in \partial M$ . Then

$$||r_x - r_y||_{C(\partial M)} \ge |r_x(z) - r_y(z)| = d(x, y)$$

Then (M, d) is isometric to  $(R(M), \|\cdot\|_{\infty})$ .

### **Lemma 2.15** The set R(M) is homeomorphic to (M, g). **Proof.**

Recall the following simple result from topology:

Assume that *X* and *Y* are Hausdorff spaces, *X* is compact and  $F: X \rightarrow Y$  is a continuous, bijective map from *X* to *Y*. Then *F* is a homeomorphism. Clearly,  $R: M \to R(M)$  is surjective and continuous. Next we prove that it is one-to-one. Assume that  $r_x(\cdot) = r_y(\cdot)$ . Denote by  $z_0$  any point where

$$d(x, \partial M) = \min_{z \in \partial M} r_x(z) = r_x(z_0) \quad \text{or}$$
$$d(y, \partial M) = \min_{z \in \partial M} r_y(z) = r_y(z_0).$$

Then  $z_0$  is a nearest boundary point to x implying that the shortest geodesic from  $z_0$  to x is normal to  $\partial M$ . The same is true for y with the same point  $z_0$ . Thus  $x = \gamma_{z_0}(s) = y$  for  $s = d(x, z_0)$ .

#### **Boundary normal coordinates.**

Consider a normal geodesic  $\gamma_z(s)$  starting from z. For small s,

$$d(\gamma_z(s), \partial M) = s, \tag{2}$$

and z is the unique nearest point to  $\gamma_z(s)$  on  $\partial M$ . Let  $\tau(z)$  be the largest value for which (2) is valid. Then for  $s > \tau(z)$ ,

$$d(\gamma_z(s), \partial M) < s,$$

and z is no more the nearest boundary point.

 $\tau(z) \in C(\partial M)$  is the cut locus distance function. The cut locus is

$$\omega = \{ x_z : z \in \partial M, \ x_z = \gamma_z(\tau(z)) \}.$$

In domain  $M \setminus \omega$  we can use the coordinates

$$x \mapsto (z(x), t(x)),$$

where  $z(x) \in \partial M$  is the unique nearest point to x and  $t(x) = d(x, \partial M)$ .

We will now use boundary normal coordinates to introduce a differential structure and metric tensor,  $g_R$ , on R(M) to have an isometry

$$R: (M,g) \to (R(M),g_R).$$

We will concentrate mainly on doing so for  $R(M) \setminus R(\omega)$ . (For the general case, see [KKL]) First, observe that we can identify those  $r = r_x \in R(M)$  with  $x \in M \setminus \omega$ .

Indeed,  $r = r_x$  with  $x = \gamma_z(s)$ ,  $s < \tau(z)$  if and only if

*i.*  $r(\cdot)$  has a unique global minimum at some point  $z \in \partial M$ .

ii. there is  $\widetilde{r} \in R(M)$  having a unique global minimum at the same z and  $r(z) < \widetilde{r}(z)$ .

A differential structure on  $R(M \setminus \omega)$  can be defined by introducing coordinates near each  $r^0 \in R(M \setminus \omega)$ . In a sufficiently small neighbourhood  $V \subset R(M)$  of  $r^0$  the coordinates

$$r \mapsto (Y(r), T(r)) = (y(\operatorname*{argmin}_{z \in \partial M} r), \min_{z \in \partial M} r)$$

are well defined. The

$$x \mapsto (Y(r_x), T(r_x))$$

coincides with the boundary normal coordinates  $x \mapsto (y(x), t(x))$  on (M, g). These coordinate determine the differential structure on  $R(M \setminus \omega)$ .

#### Construction of the metric $g_R$ on R(M).

Let  $r^0 \in R(M \setminus \omega)$ ,  $V \subset R(M)$  be its neighbourhood, and  $Y: V \to U \subset \mathbb{R}^m$  be local coordinates,  $Y(r^0) = 0$ 

For  $z \in \partial M$  we define an evaluation function

$$K_z: V \to \mathbb{R}, \quad K_z(r) = r(z).$$

The function  $E_z = K_z \circ Y^{-1} : U \to \mathbb{R}$  satisfies

$$E_z(y) := d(z, Y^{-1}(y)), \quad y \in U.$$

Consider the function  $E_z(y)$  as a function of y with a fixed z. The differential  $dE_z$  at point 0 is a covector in  $T_0^*U$ . Since the gradient of a distance function has length one, we see that

$$||dE_z||_{g_R}^2 := (g_R)^{jk} \frac{\partial E_z}{\partial y^j} \frac{\partial E_z}{\partial y^k} = 1, \quad j, k = 1, \dots, m.$$

Varying  $z \in \partial M$  we obtain a set of covectors  $dE_z(0)$  in the unit ball of  $(T_0^*U, g_R)$  which contains an open set.

This determines uniquely the tensor  $g_R$ .

Hence we have proven

**Theorem 5** The boundary data  $(\partial M, \Lambda)$  determine the manifold (M, g) upto isometry.

Also the potential q(x) of the operator  $-\Delta_g + q$  can be uniquely determined.

## 2.16 New results: Time reversal

On formal level, the the previous algorithm is based on the following task: Let f be given. Can we find h such that

$$u^{h}(x,T) = \chi_{M(\Gamma,\tau)}(x)u^{f}(x,T).$$

This is equivalent of the minimization of

$$||u^{f}(T) - u^{h}(T)||_{L^{2}(M)} : \quad h \in C_{0}^{\infty}(\Gamma \times [0, \tau]).$$



Generally, the minimization problem has no solution and is ill-posed. We consider the regularized minimization problem

$$\min_{h \in L^2(\partial M \times [0,2T])} F(h,\alpha)$$

where  $\alpha \in (0, 1)$  and

$$F(h,\alpha) = \langle K(Ph-f), Ph-f \rangle_{L^2(\partial M \times [0,2T], dS_g)} + \alpha \|h\|_{L^2}^2.$$

Let us recall the Blagovestchenskii identity

$$\int_{M} u^{f}(x,T)u^{h}(x,T) dV_{\mu}(x)$$

$$= \int_{[0,2T]^{2}} \int_{\partial M} J(t,s)[f(t)(\Lambda_{2T}h)(s) - (\Lambda_{2T}f)(t)h(s)]dS_{g}dtds$$

$$= \int_{\partial M \times [0,2T]} (Kf)(x,t)h(x,t) dS_{g}(x)dt,$$

where  $J(t,s) = \frac{1}{2}\chi_L(s,t)$  and

$$L = \{ (s, t) \in \mathbb{R}_+ \times \mathbb{R}_+ : t + s \le 2T, s > t \}.$$

Here

$$K = R_{2T}\Lambda_{2T}R_{2T}J - J\Lambda_{2T},$$

where

$$Rf(x,t) = f(x,2T-t),$$

is the time reversal operator and

$$Jf(x,t) = \frac{1}{2} \int_0^{\min(2T-t,t)} f(x,s) ds,$$

is the time filter. Note that

 $\Lambda_{2T}^* = R_{2T}\Lambda_{2T}R_{2T} \quad \text{as} \quad G(x, x', t' - t) = G(x', x, -(t) - (-t')).$ 

We also use the restriction operator

$$P_B f(x,t) = \chi_B(x,t)u(x,t),$$

The processed time reversal iteration is

$$F := \frac{1}{\omega} P(R\Lambda_{2T}RJ - J\Lambda_{2T})f,$$
  

$$a_n := \Lambda_{2T}(h_n),$$
  

$$b_n := \Lambda_{2T}(RJh_n),$$
  

$$h_{n+1} := (1 - \frac{\alpha}{\omega})h_n - \frac{1}{\omega}(PRb_n - PJa_n) + F,$$

where  $f \in L^2(\partial M \times [0, 2T])$  and  $\alpha, \omega > 0$  are parameters. Iteration starts at  $h_0 = 0$ . **Theorem 6 (Bingham-Kurylev-L.-Siltanen 2007)** Let  $\Gamma_1 \subset \partial M$ ,  $0 \leq T_1 \leq T$ , and  $B = \Gamma_1 \times [T - T_1, T]$ . Let  $f \in L^2(\partial M \times \mathbb{R}_+)$  and  $h_n = h_n(\alpha)$  be defined by the processed time reversal iteration. Then

$$h(\alpha) = \lim_{n \to \infty} h_n(\alpha)$$

and the limits satisfy in  $L^2(M)$ 

$$\lim_{\alpha \to 0} u^{h(\alpha)}(x,T) = \chi_{M(\Gamma_1,T_1)}(x) u^f(x,T).$$



#### **Proof.** The minimization problem

 $\min_{h \in L^2(\partial M \times [0,2T])} F(h,\alpha)$ 

with  $\alpha \in (0,1)$  and

$$F(h,\alpha) = \langle K(Ph-f), Ph-f \rangle_{L^2(\partial M \times [0,2T], dS_g)} + \alpha \|h\|_{L^2}^2$$

leads to a linear equation

$$(PKP + \alpha)h = PKf.$$

This can be solved using iteration.

**Corollary 2.17** Assume we are given the boundary  $\partial M$ and the response operator  $\Lambda$ . Then using the the processed time reversal iteration we can find constructively the manifold (M, g) upto an isometry and on it the operator Auniquely.


Let  $x = \gamma_{z,\nu}(s)$ .

The distance dist (x, z) is the infimum of all  $\tau$  that satisfy the condition

(A) The set

$$(M(z,s) \cap M(y,\tau)) \setminus M(\partial M, s-\epsilon)$$

is non-empty for all  $\epsilon > 0$ .

## References:

- 1. M. Belishev: An approach to multidimensional inverse problems for the wave equation. (Russian) *Dokl. Akad. Nauk SSSR* 297 (1987), no. 3, 524–527.
- 2. M. Belishev, Y. Kurylev: A nonstationary inverse problem for the multidimensional wave equation 'in the large." (Russian) *Zap. Nauchn. Sem. LOMI* **165** (1987), 21–30.
- 3. K. Bingham, Y. Kurylev, M. Lassas, S. Siltanen: Time reversal methods in unknown medium and inverse problems, preprint, 2007.

- 4. A. Katchalov, Y. Kurylev, M. Lassas: *Inverse Boundary Spectral Problems*, Monographs and Surveys in Pure and Applied Mathematics 123, Chapman Hall/CRC-press, 2001, xi+290 pp.
- A. Katchalov, Y. Kurylev, M. Lassas, N. Mandache: Equivalence of time-domain inverse problems and boundary spectral problem, Inverse problems. 20 (2004), 419-436
- A. Katchalov, Y. Kurylev, M. Lassas Energy measurements and equivalence of boundary data for inverse problems on non-compact manifolds. IMA volumes in Mathematics and Applications (Springer Verlag) 'Geometric Methods in Inverse Problems and PDE Control" Ed. C. Croke, I. Lasiecka, G. Uhlmann, M. Vogelius, 2004, pp. 183-214.

- Y. Kurylev, M. Lassas: Multidimensional Gel'fand Inverse Boundary Spectral Problem: Uniqueness and Stability, CUBO Mathematical Journal 8 (2006), no. 1, 41–59
- 8. Y. Kurylev, M. Lassas, E. Somersalo: Maxwell's equations with a polarization independent wave velocity: Direct and inverse problems, *Journal de Mathmatique Pures et Appliques* 86 (2006), no. 3, 237-270.
- 9. Y. Kurylev, M. Lassas: Inverse Problems and Index Formulae for Dirac Operators, preprint 2006.

 J. Sylvester, G. Uhlmann: A global uniqueness theorem for an inverse boundary value problem. *Ann. of Math.* (2) 125 (1987), no. 1, 153–169.